

Bending of a piecewise homogeneous plate with a circular interfacial materials separation zone and radial crack considering the strip contact of its edges

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Abstract. The work presents a solution to the bending problem of an infinite piecewise homogeneous isotropic plate with an elastic circular washer and a radial through straight crack. It was assumed that under the action of an external loads at infinity, the edges of the crack are smoothly contacted on area of constant width (strip contact) on the upper base of the plate. The solution of the problem is built using the methods of the theory of functions of a complex variable and complex potentials and is reduced to a system of singular integral equations, which is numerically solved using the method of mechanical quadrature. A numerical analysis of the problem is conducted and graphic dependencies of contact force, coefficients of intensity of moments and forces at various parameters of the problem were constructed.

Keywords: plate, crack, closure, plane problem, bending, linear conjugation, singular equations, contact force.

1. Introduction

Lamellar structural elements are widely used in various fields of engineering. They may contain finite inclusions of a technological nature and during operation, cracks may appear in such elements, abruptly reducing the permissible load level that the structure can withstand. It is clear from the physical considerations that, under the conditions of plate bending, the crack edges will be in contact. Therefore, from scientific and technical task of mechanics it is relevant to develop methods that determine the stress-strain state of plates with cracks which edges are in contact during bending. One of the main tasks in assessing the strength of a structural element with cracks is to determine the coefficients of the intensity of forces and moments, which are used to characterize the stress-strain state around the crack peaks, as well as to determine the critical load that can be applied to the plate.

Research on the stress-strain state of plates with cracks during bending using the Kirchhoff-Love theory is given in [1-4]. In these works, it was assumed that the edges of the crack are in contact along the entire length on the upper or lower base of the plate. Kalthoff J. F., Shatsky I. P., Bürgel A. [5] experimentally investigated the contact of crack edges during the bending of a plate and confirmed the presence of a membrane component of the stress field near the crack, which is a direct consequence of the contact interaction of the crack edges during the bending of the plate.

In the work of J. P. Dempsey, I. I. Shekhtman, and L. I. Slepyan [6], the problem of bending a

plate with a through crack, where the edges of the crack are in contact over the height of the plate, was the first time analytically solved. The bending of a circular plate with a system of through straight cracks, the edges of which are in contact over the height of the plate, is considered in [7]. The stress-strain state of a plate with a crack, the edges of which are in contact under the action of external load, and on the continuation of which plastic zones were formed, was investigated in works [8, 9]. The study of the stress-strain state of bodies with interphase cracks with contact zones was carried out in works [10, 11]. The influence of the contact of the edges of a thermoisolated interfacial crack under thermomechanical loading on the stress-strain state was studied in [12]. The stress-strain state of a beam with a crack, the edges of which are in contact and on the continuation of which a plastic zone was formed, was investigated in [13]. Therefore, problems of bending plates with a system of cracks, the edges of which are in contact, are relevant. The work formulates a model of the smooth contact of the crack edges on area of constant width for the bending problem of a piecewise homogeneous plate. This model makes it possible to estimate more accurately the influence of a crack on the stress-strain state of the plate and on its strength.

2. Problem statement

Consider an infinite piecewise homogeneous isotropic plate of thickness $2h$ with a circular interfacial materials separation zone of radius R and a radial crack of length $2l$, which is free from external loading. It is assumed that under the action of bending moments at infinity (Fig. 1(a)) the edges of the crack come into smooth contact along the region of constant width (strip contact) on the upper base of the plate (Fig. 1(b)).

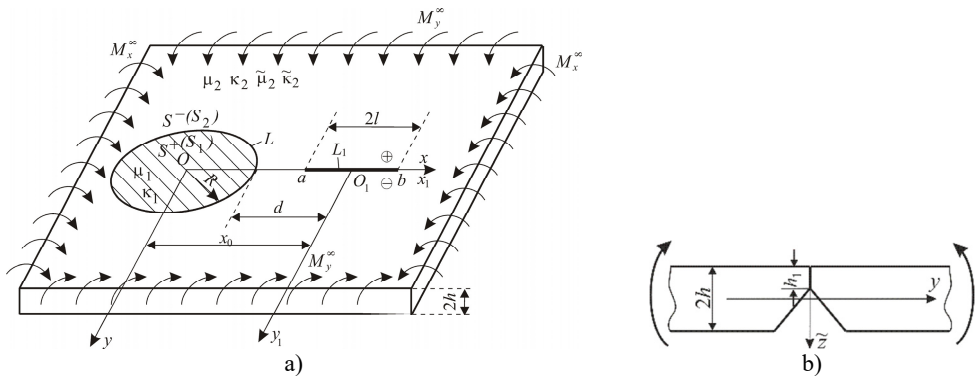


Fig. 1. a) The scheme of the plate loading and crack placement and b) the scheme of the crack edge contact

In the middle plane of the plate, we choose a Cartesian coordinate system $Oxyz$, directing the axis Oz perpendicular to it, and the origin of the coordinates O coincides with the center of the circular washer, and the crack located along the Ox . Additionally, in this plane, we introduce a polar coordinate system r and θ with the pole at the point O and the polar axis along the Ox . We assume that the center of the crack has the coordinate x_0 , where $x_0 > R + l$. We associate a Cartesian coordinate system $O_1x_1y_1$ with the crack. Quantities related to the washer will be denoted with index 1, and those related to the matrix with index 2. The points of the plane Oxy , coinciding with the crack ends will be denoted by a and b , the area in the middle of the circular tough washer – as $S^+(S_1)$, outside – as $S^-(S_2)$, the line where the crack is located – as L_1 , and the separation boundary of materials – by L . The distributed bending moments at infinity are denoted by M_x^∞ and M_y^∞ (Fig. 1(a)).

Due to the contact of the crack edges, the solution to the problem will be sought in the form of solutions of two interrelated problems: the problem of plate bending (classical plate bending theory) and the plane problem of elasticity theory under the following boundary conditions:

$$\sigma_{y_1 y_1}^\pm = \frac{-N}{(2h)}, \quad \sigma_{x_1 y_1}^\pm = 0, \quad P^\pm = 0, \quad \partial_{x_1}[u_{\Pi 1}] + \alpha h[\partial_{x_1 y_1}^2 w] = 0, \quad (1)$$

$$M_{y_1}^\pm = \beta h N, \quad x_1 \in L_1, \\ P_{r1} = P_{r2}, \quad M_{r1} = M_{r2}, \quad u_{r1} = u_{r2}, \quad u_{\theta 1} = u_{\theta 2}, \quad w_1 = w_2, \quad \partial_r w_1 = \partial_r w_2, \quad x \in L, \quad (2)$$

where $\alpha = 0,5 \cdot (1 + (1 - \gamma)^2)$, $\beta = 1 - \gamma/3$, $\gamma = h_1/h$, N – contact force between the edges of the crack; $\sigma_{x_1 y_1}$, $\sigma_{y_1 y_1}$ – components of the stress tensor, u_{rj} and $u_{\theta j}$, $u_{\Pi 1}$ – components of the displacement vector in the plane problem; here and henceforth, the index j takes two values: 1 and 2; w_j – deflection of the plate, M_{rj} , M_{y_1} – bending moments, P_{y_1} , P_{rj} – shearing force generalized in the sense of Kirchhoff; $[f] = f^+ - f^-$, the “+” and “-” signs indicate the boundary values of the function when the point of the plane approaches the crack at $y_1 \rightarrow \pm 0$, $\partial_x f = \partial f / \partial x$.

3. Solution

Using the complex potentials of the plane problem and the plate bending problem [14], satisfying the boundary conditions Eqs. (1)-(2), the solution of the problem is reduced to linear conjugation problems, based on which a system of singular integral equations is obtained for the unknown displacement jumps $G_j(\eta)$ and rotation angle jumps $Y_j(\eta)$ on the crack edges, which takes the form:

$$i \cdot \int_{-1}^1 \left\{ Y_1(\eta) \cdot K(\eta, \xi) + \overline{Y_1(\eta)} \cdot L(\eta, \xi) \right\} d\eta = i\tilde{c}' + \tilde{m} \cdot \frac{\beta h N}{M_y^\infty} + P(\xi), \quad \xi \in [-1, 1], \quad (3)$$

$$\int_{-1}^1 \left\{ G_1(\eta) \cdot R(\eta, \xi) + \overline{G_1(\eta)} \cdot S(\eta, \xi) \right\} d\eta = \frac{-0,5\pi \cdot h N}{M_y^\infty}, \quad \xi \in [-1, 1], \quad (4)$$

where:

$$K(\eta, \xi) = -\frac{1}{\pi} \left\{ \frac{\tilde{\kappa}_2}{\eta - \xi} + \frac{\lambda}{2} \left\{ \tilde{g} \tilde{A}_4 \left(1 + \frac{1}{X^2} \right) + (\tilde{a} - \tilde{A}_5 \tilde{A}_{13}) \frac{1}{2X^2} \right\} \frac{1}{T} \right. \\ \left. + \frac{1}{TX - 1} \left\{ \tilde{g} \tilde{A}_4 \left(\frac{\tilde{\kappa}_2}{X} - X - \frac{3}{X^3} \right) - \frac{\tilde{A}_3}{\tilde{g}X} \right\} \right. \\ \left. + \frac{\tilde{g} \tilde{A}_4}{(TX - 1)^2} \left(X + \frac{4}{X} - \frac{5}{X^3} \right) - \frac{2\tilde{g} \tilde{A}_4}{X} \left(X - \frac{1}{X} \right) \frac{1}{(TX - 1)^3} \right\}, \\ L(\eta, \xi) = -\frac{\lambda}{2\pi} \left\{ \frac{1}{T} \left(\tilde{\kappa}_2 \tilde{g} \tilde{A}_4 - \frac{1}{X^2} \left(\tilde{A}_4 - \frac{\tilde{a}}{2} + \frac{\tilde{A}_5 \tilde{A}_{13}}{2} \right) \right) + \frac{\tilde{\kappa}_2 \tilde{g} \tilde{A}_4}{TX - 1} \left(\frac{3}{X^3} - X - \frac{2}{X} \right) \right. \\ \left. + \frac{\tilde{\kappa}_2 \tilde{g} \tilde{A}_4}{(TX - 1)^2} \left(X + \frac{1}{X^3} - \frac{2}{X} \right) \right\}, \\ P(\xi) = -\tilde{m} + \tilde{g} \tilde{A}_4 \frac{B'}{X^2} \left(\tilde{\kappa}_2 + 1 - \frac{3}{X^2} \right) - 2\tilde{A}_{13} \frac{A}{X^2}, \\ R(\eta, \xi) = \frac{1}{\eta - \xi} - \frac{\lambda g A_4}{2(TX - 1)} \left\{ \frac{1}{X} - \left(\frac{1}{X^2} + 1 \right) \left\{ X + \frac{1}{X} - \left(X - \frac{1}{X} \right) \frac{1}{TX - 1} \right\} \right. \\ \left. + \left(X - \frac{1}{X} \right) \left\{ \frac{4}{X^2} \frac{1}{TX - 1} + \frac{2}{X^2} - \frac{2}{X} \left(X - \frac{1}{X} \right) \frac{1}{(TX - 1)^2} \right\} \right\} + \frac{\lambda g^{-1} A_3}{2X(TX - 1)} + \frac{\lambda A_4}{2T} \left(\frac{B_8}{X^2} - g \right), \\ S(\eta, \xi) = -\frac{\lambda}{2} \left\{ \frac{1}{T} \left(\frac{B_9}{X^2} + g A_4 \right) - \frac{g A_4}{TX - 1} \left\{ X - \frac{1}{X^3} - \left(X - \frac{1}{X} \right) \frac{1}{TX - 1} + \frac{1}{X^2} \left(\frac{1}{TX - 1} + 2 \right) \right\} \right\},$$

$$\begin{aligned}
 Y_1(t) &= \frac{Q_2 Q(t)}{M_y^\infty} = Y_{11}(t) + iY_{12}(t), \quad G_1(t) = -ih^2(M_y^\infty)^{-1}g'(l\eta) = G_{11}(\eta) + iG_{12}(\eta), \\
 Q_j &= E_j h^3, \quad D_j = 2Q_j \left(3(1 - \nu_j^2)\right)^{-1}, \quad \tilde{m} = -\left(\tilde{D}_2(1 - \nu_2)\right)^{-1}, \quad \tilde{D}_2 = 2\left(3(1 - \nu_2^2)\right)^{-1}, \\
 \tilde{\kappa}_j &= \frac{(3 + \nu_j)}{(1 - \nu_j)}, \quad B = -\frac{1 - \rho}{2}\tilde{m}, \quad \mu_j = \frac{E_j}{\left(2(1 + \nu_j)\right)}, \quad \kappa_j = \frac{(3 - \nu_j)}{(1 + \nu_j)}, \\
 \tilde{\mu}_j &= \frac{1}{D_j(1 - \nu_j)}, \quad A = -\frac{\rho + 1}{4\tilde{D}_2(1 + \nu_2)}, \quad \rho = \frac{M_x^\infty}{M_y^\infty}, \\
 \tilde{B}'_1 &= 0, \quad \tilde{B}'_0 = -\tilde{A}'_0, \quad \tilde{A}_j = \tilde{\mu}_j + \tilde{\mu}_{3-j}\tilde{\kappa}_j, \quad \tilde{g} = -\tilde{A}_1\tilde{A}_2^{-1}, \\
 \tilde{A}A_3 &= \tilde{\mu}_1\tilde{\kappa}_2 - \tilde{\mu}_2\tilde{\kappa}_1, \quad \tilde{A}A_4 = \tilde{\mu}_2 - \tilde{\mu}_1, \quad \tilde{A}_4 = \tilde{A}A_4\tilde{A}_1^{-1}, \quad \tilde{A}_3 = \tilde{A}A_3\tilde{A}_2^{-1}, \\
 \tilde{A}_5 &= \tilde{\mu}_2(1 + \tilde{\kappa}_1)\tilde{A}_1^{-1}, \quad \tilde{A}_{12} = \tilde{A}_4 - \tilde{g}^{-1}\tilde{A}_3, \quad \tilde{a} = \tilde{A}_3\tilde{g}^{-1} + \tilde{A}_4, \quad \tilde{A}_{13} = \tilde{A}_{12}/(1 - \tilde{A}_4), \\
 B_8 &= A_5(g^{-1}A_3 - 1)A_6, \quad B_9 = A_4 + A_5(A_4^2 - g^{-1}A_3)A_6, \quad B_5 = \mu_2(1 + \kappa_2)A_1^{-1}, \\
 \tilde{A}_3 &= \mu_1\kappa_2 - \mu_2\kappa_1, \quad \tilde{A}\tilde{A}_4 = \mu_2 - \mu_1, \quad g = -A_a A_2^{-1}, \quad A_3 = \tilde{A}\tilde{A}_3\tilde{A}_2^{-1}, \quad A_4 = \tilde{A}\tilde{A}_4\tilde{A}_1^{-1}, \\
 \tilde{A}A_j &= \mu_j + \mu_{3-j}\kappa_j, \quad A_6 = (1 - A_4^2)^{-1}, \quad B_8 = A_5(g^{-1}A_3 - 1)A_6, \\
 B_9 &= A_4 + A_5(A_4^2 - g^{-1}A_3)A_6, \quad B_5 = \mu_2(1 + \kappa_2)A_1^{-1}, \quad T = \tilde{X}_0 + \lambda\eta, \quad X = \tilde{X}_0 + \lambda\xi,
 \end{aligned}$$

where E_j – Young’s modulus, ν_j – Poisson’s ratios; $Y_{11}(t)$, $Y_{12}(t)$, $G_{11}(t)$ and $G_{12}(t)$ – real functions.

To the Eqs. (3), (4) we add additional conditions:

$$\int_{-1}^1 Y_1(\eta)d\eta = 0, \quad \int_{-1}^1 \eta Y_{11}(\eta)d\eta = 0, \quad \int_{-1}^1 G_1(\eta)d\eta = 0, \quad (5)$$

which express respectively single-valued rotational angles and deflection of the plate when bypassing the contour of each crack.

From the boundary Eq. (1), we obtain:

$$G_{11}(\eta) + \frac{\alpha(1 + \tilde{\kappa}_2)}{(1 + \kappa_2)(1 + \nu_2)}Y_{12}(\eta) = 0, \quad \eta \in [-1, 1]. \quad (6)$$

Thus, to determine the unknown functions $Y_1(\eta)$ and $G_1(\eta)$ we have a system of integral Eqs. (3)-(6).

4. Numerical analysis and conclusions

A numerical analysis of the problem was carried out, which is presented in Fig. 2-4 at $\nu_1 = \nu_2 = 0.3$, $\gamma = 0.13$. It should be noted, that the consolidated coefficients of the intensity of forces $K_N^* = K_N/(M_y^\infty\sqrt{l})$ and moments $K_M^* = K_M/(M_y^\infty\sqrt{l})$ are related by the dependence $K_N^*/K_M^* = 3\alpha(1 + \nu_2)/(3 + \nu_2)$, therefore, we do not provide graphical dependencies for K_N^* .

Fig. 2a illustrates graphical dependence of the consolidated contact force between the edges of the crack $N^* = hN/M_y^\infty$ on the consolidated coordinate $\xi = x_1/R$ for $\tilde{n} = E_1/E_2 = 0.1$, $\rho = M_x^\infty/M_y^\infty = 1$, $R = l$. Curve 1 is obtained at $\varepsilon = \lambda = 1/1.1$ ($\varepsilon = R/d$, $\lambda = l/d$), curve 2 – at $\varepsilon = \lambda = 1/1.2$, curve 3 – at $\varepsilon = \lambda = 1/1.5$, curve 4 – at $\varepsilon = \lambda = 1/2$, curve 5 – at $\varepsilon = \lambda = 0$. As we can see from this figure, as the crack approaches the circular boundary, the magnitude of the contact force N^* between the edges of the crack increases compared to an isolated crack, and for the peak closer to the washer a ($\xi = -1$) it is larger than farther from the hole b ($\xi = 1$).

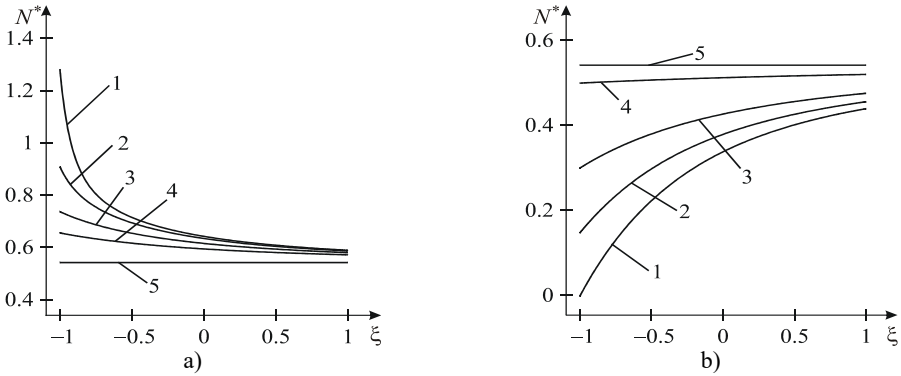


Fig. 2. Graphical dependence of the consolidated contact force between the edges of the crack

Fig. 3 and Fig. 4 shows the graphical dependence for coefficients of the intensity of moments at vertices *a* and *b* on $\varepsilon = R/d$ at $\tilde{n} = E_1/E_2 = 0.1$ and $\tilde{n} = E_1/E_2 = 10$ respectively, when $\rho = M_x^\infty/M_y^\infty = 1$. Solid lines represent the curves for coefficients of the intensity of moments, obtained, taking into account the contact of the crack edges, dashed lines – without considering the contact between the crack edges. Curves 1 are constructed at $\lambda = l/d = 0.2$, curves 2 – at $\lambda = 0.5$, curves 3 – at $\lambda = 0.8$. As can be seen from Fig. 4 when $\tilde{n} = E_1/E_2 = 0.1$ coefficients of the intensity of moments increases with the increase of the relative radius $\varepsilon = R/d$ of the circular interfacial materials separation zone, but at $\tilde{n} = E_1/E_2 = 10$ (Fig. 4) coefficients of the intensity of moments with the same parameters of the problem decrease with the increase of $\varepsilon = R/d$.

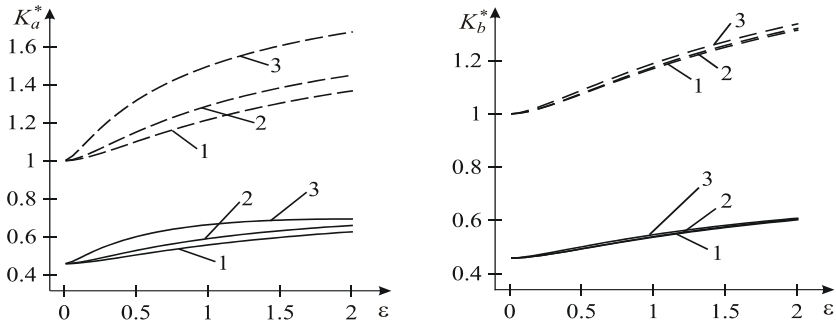


Fig. 3. Graphical dependence of the consolidated coefficients of the intensity of moments on ε for $\tilde{n} = 0.1$

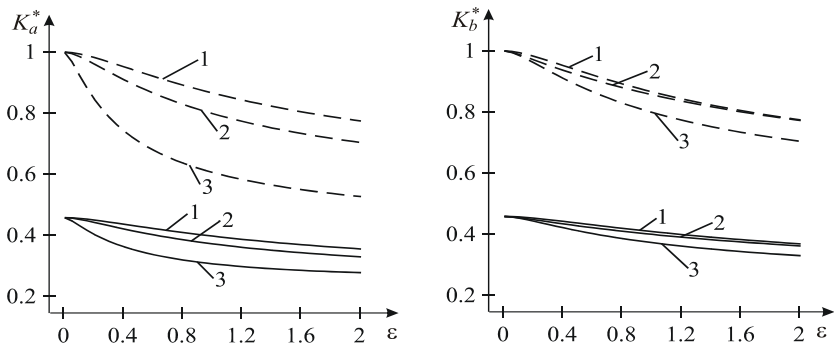


Fig. 4. Graphical dependence of the consolidated coefficients of the intensity of moments on ε for $\tilde{n} = 10$

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Data availability

The datasets generated during and/or analyzed during the current study are available from the corresponding author on reasonable request.

Conflict of interest

The authors declare that they have no conflict of interest.

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