Dynamic characteristics analysis of a coaxial rotor system with shared bearing bore structure

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Abstract. The nonlinear dynamic model of a shared bearing bore-coaxial rotor system was established with nonlinear forces of the rolling bearings and squeeze film dampers considered, Cross coupling frequency was found in cross coupling vibration simulation. Numerical studies were conducted to reveal the mechanism of the coupling vibration for the shared bearing bore-coaxial rotor system. The coupling vibration characteristics and influencing factors of the shared bearing bore-coaxial rotor system were studied. The results show that the coupling vibration is featured by cross-excitation phenomenon and orbit which shows a "circle in circle" pattern. Besides, the angular stiffness of the shared bearing bore has great influences on the coupling vibration. However, the radial stiffness has almost no effect on the coupling vibration.

Keywords: shared bearing bore structure, nonlinear vibration, coupling vibration, coaxial rotor system.

1. Introduction

The shared bearing bore structure has been widely used in high power-to-weight ratio turboshaft engines. It is one of the most important technologies to reduce engine weight, increase power-to-weight ratio and turbine efficiency [1]. The shared bearing bore structure supports both the gas generator and the power turbine rotor, thus to reduce the number of bearing cavity and supports and then the weight of engine [2]. However, vibration of the gas generator would be transmitted to the power turbine rotor and vice versa. And coupling vibration would increase the opportunities of resonance for both rotors [3]. Besides, the power turbine is a typical overhung high-speed flexible rotor system, which is highly sensitive to external excitations, such as the unbalance load and the vibration transmitted by the shared bearing bore. Therefore, it is important to study the coupling dynamics of the shared bearing bore-dual rotor system.

At present, there are very few studies about the shared bearing bore at home and abroad. Most studies of coupling dynamics for dual rotor system with inter-shaft bearing focus on cross excitation phenomena. Gupta et al. [4-5] established the dynamic model of a rolling bearing-dual rotor system by the transfer matrix method. The influences of rolling bearing, moments of inertia and gyroscopic effect were studied and verified. Guskov et al. [6-9] conducted numerical and experimental researches on a dual rotor system to study the coupled unbalance response characteristics. Bonello et al. [10-11] developed a modeling method in frequency domain to study the nonlinear dynamic characteristics of the squeeze film damper (SFD)-dual rotor system. Luo Guihuo et al. [12-13] studied the nonlinear dynamics of the co-and counter-rotating dual rotor system with inter-shaft bearing and SFD. Also, the influences of radial clearance of the inter-shaft bearing were studied. Chen Guo et al. [14-16] established the dynamic model of a casing-rolling bearing-rotor system. The coupling vibration between the rotors were studied. But the influences of the supporting structures were neglected.

The shared bearing bore-coaxial rotor system differs from the rolling bearing-dual rotor system in that the dynamic characteristics of the supporting structure itself have great impact on the rotor's

response. Thus, to promote detailed studies on the dynamic characteristics of the shared bearing bore-coaxial rotor system. the remainder of this paper is organized as follows: the modeling method is presented in Section 2; the influences of radial and the angular stiffness are studied in Section 3; the main conclusions are drawn and discussed in Section 4.

2. Modeling

2.1. The coaxial rotor system

The shared bearing bore coaxial rotor system is shown in Fig. 1. The rotor system consists of the power turbine rotor, the gas generator rotor, the shared bearing bore and squeeze film dampers. The power turbine rotor is supported by bearing 1, 2, 5 and 6. The gas generator is supported by bearing 3 and 4. Bearing 4, 5 and 6 are connected to the outer case by the shared bearing bore. Bearing 1 and 5 adopts the rolling bearing+squirrel cage scheme while the rest supports adopt rolling bearing+squirrel cage+squeeze film damper scheme. Relevant parameters of rotor system are listed in Table 1 to Table 4. In Table 1, k_{xx} and k_{xz} indicate the radial and angular stiffness, respectively. In Table 2, L represents the length of the oil film, R is radius the oil film, and C is the radial clearances. In Table 3, I_p indicates the polar inertia of the disk.

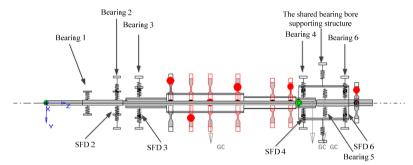


Fig. 1. Schematic diagram of the Shared bearing bore-coaxial rotor

Table 1. Stiffness of squirrel cages

k_{xx} / (N/m)	k_{xz} / (N·m/rad)
3.77×10^{6}	0
2.35×10 ⁵	0
1.13×10 ⁶	0
7.10×10 ⁵	0
4.13×10 ⁶	0
2.30×10 ⁵	0
	3.77×10 ⁶ 2.35×10 ⁵ 1.13×10 ⁶ 7.10×10 ⁵ 4.13×10 ⁶

Table 2. Parameters of SFD

	Support number	L / mm	R / mm	<i>C</i> / mm	Dynamic viscosity / (Pa·s)	
Power turbine rotor	2	19.0	23.0	0.05	1.0752×10 ⁻²	
	6	15.5	26.5	0.05		
Gas generator rotor	3	9.5	25.5	0.05	1.0/32×10 2	
	4	17.0	26.5	0.05		

2.2. Equations of motion

The Timoshenko beam element is adopted to model the shaft and the shared bearing bore. Each beam element consists of 2 nodes with 4 degrees of freedom per node. Translational degrees of freedom in x, y direction u, v and rotational degrees of freedom around x, y axis, respectively, as shown in Fig. 2.

		i abie 3. Para	meters of a	ISKS		
	Section	Starting point /	Mass /	I_p /	Unbalance /	Initial
	number	mm	kg	(kg·m ²)	(g·cm)	phase
Gas generator rotor	1	467	4.66	0.035890	2	0
	2	543	5.26	0.044880	2	0
	3	615	2.01	0.008250	1	0
	4	718.5	0.53	0.001110	0	0
	5	849	2.49	0.009724	2	0
	6	916.5	1.48	0.003764	2	0
Power turbine rotor	7	1164	5 16	0.072610	5	0

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The node displacement vector of the element can be written as:

$$\mathbf{q}_e = [u_{e1} \quad v_{e1} \quad \theta_{e1} \quad \varphi_{e1} \quad u_{e2} \quad v_{e2} \quad \theta_{e2} \quad \varphi_{e2}]^T. \tag{1}$$

The equations of motion of the beam element:

$$\mathbf{M}_{e}\ddot{\mathbf{q}}_{e} + \omega \mathbf{G}_{e}\dot{\mathbf{q}}_{e} + \mathbf{K}_{e}\mathbf{q}_{e} = \mathbf{Q}_{e},\tag{2}$$

where, \mathbf{Q}_e represents the external force load of the beam element; \mathbf{M}_e represents the inertia matrix; and G_e represents the gyroscopic matrix; K_e represents the stiffness matrix; ω is the rotational speed.

The equations of motion of the disc can be written as:

$$\mathbf{M}_{d}\ddot{\mathbf{q}}_{d} + \omega \mathbf{G}_{d}\dot{\mathbf{q}}_{d} = \mathbf{Q}_{d},\tag{3}$$

where \mathbf{Q}_d - the generalized external force vector, \mathbf{M}_d - the inertia matrix, \mathbf{G}_d - the gyroscopic

The expressions for \mathbf{M}_d and \mathbf{G}_d are shown in Eqs. (4) and (5). Where: m_d indicates the mass of the disc; I_d and I_p indicate the diametral and polar moment inertia of the disc, respectively:

A schematic diagram of the supporting structure is shown in Fig. 4. Thus the supporting force consists of three parts, namely the nonlinear force f_b of the rolling bearing, the nonlinear force of the squeeze film damper f_s and the elastic supporting force f_{sq} . The elastic supporting force is:

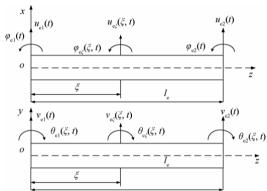
$$f_{Sq} = k_S q_{S'} \tag{6}$$

where, k_s is the stiffness of the elastic support, and q_s is the displacement of the journal.

Based on the short-bearing assumptions and the Reynolds boundary conditions, nonlinear forces of the squeeze film damper can be expressed as:

$$\begin{cases} f_{sx} = -\frac{\mu R L^3}{c^2 (x^2 + y^2)^{1/2}} \left[x (\dot{\varepsilon} I_2 + \varepsilon \dot{\psi} I_1) + y (\dot{\varepsilon} I_1 + \varepsilon \dot{\psi} I_3) \right] \\ f_{sy} = -\frac{\mu R L^3}{c^2 (x^2 + y^2)^{1/2}} \left[y (\dot{\varepsilon} I_2 + \varepsilon \dot{\psi} I_1) + x (\dot{\varepsilon} I_1 + \varepsilon \dot{\psi} I_3) \right] \end{cases}, \tag{7}$$

where x and y are the horizontal and vertical displacements of the journal. I_j (j=1,2,3) are Sommerfeld integrals. R and L are radius and length of the SFD respectively. μ and c denote the dynamic viscosity of the film and the radial clearance of the SFD. $\varepsilon = \sqrt{x^2 + y^2}/c$; $\dot{\varepsilon} = (x\dot{x} + y\dot{y})/(c\sqrt{x^2 + y^2})$; $\tan \psi = x/y$; $\dot{\psi} = (y\dot{x} - x\dot{y})/(x^2 + y^2)$.



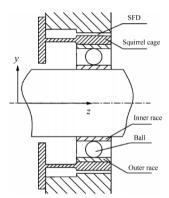


Fig. 2. Schematic diagram of beam element

Fig. 3. Schematic diagram of the support

As for the rolling ball bearing, based on pure rolling assumption and the Hertz contact theory, nonlinear force of the rolling ball bearing is:

$$\begin{cases} f_{x}^{b} = k_{n} \sum_{j=1}^{N_{b}} u_{\theta j}^{\xi} H(u_{\theta j}) \sin \theta_{j}, \\ f_{y}^{b} = k_{n} \sum_{j=1}^{N_{b}} u_{\theta j}^{\xi} H(u_{\theta j}) \cos \theta_{j}, \\ f_{y}^{b} = k_{n} \sum_{j=1}^{N_{b}} u_{\theta j}^{\xi} H(u_{\theta j}) \cos \theta_{j}, \end{cases}$$

$$(8)$$

$$\theta_{j} = \frac{2\pi(j-1)}{N_{b} + \omega_{c}t}, \quad \omega_{c} = \frac{(\omega_{in}r + \omega_{out}R)}{(R+r)}.$$

$$u_{\theta j} = (x^{ir} - x^{or}) \sin \theta_{j} + (y^{ir} - y^{or}) \cos \theta_{j} - \gamma/2,$$

The equations of motion of the beam element and discs are assembled to obtain the equation of motion of the rotor system:

$$\mathbf{M}\ddot{\mathbf{q}} + (\mathbf{C} + \omega \mathbf{G})\dot{\mathbf{q}} + \mathbf{K}\mathbf{q} = \mathbf{Q},\tag{9}$$

where, **M**, **G**, **K** represent the inertia, gyroscopic and stiffness matrix of the system, respectively. **Q** is the generalized external force vector of the system, including unbalance force, nonlinear forces of the rolling bearing and the squeeze film dampers.

Mathematically, the model is a set of nonlinear second order differential equations. The computational efficiency of which depends on the numerical methods applied and the total DOFs of the rotor system. To combine accuracy of the finite element model and computational efficiency, the fixed interface modal synthesis is applied to reduce dimension of the mathematical model and thus the computational effort. Detailed procedures can be seen in reference [12-13],

which will not be repeated here.

3. Results and discussion

3.1. Influence of the radial stiffness of the shared bearing bore structure

To investigate the influence of the shared bearing bore structure on the coupled vibration response, the radial stiffness of the shared bearing bore structure was taken at 1×10^6 N/m, 1×10^8 N/m and 1×10^{10} N/m respectively while other parameters remain unchanged. The orbit is shown in Fig. 4 and the waterfall diagram of the response of the power turbine disc is shown in Fig. 5.

It can be seen from Fig. 4 that the orbit of the turbine shows a pattern of circle in circle due to the "cross excitation". Also, the orbit plot in Fig. 4 indicates that increasing the stiffness of the shared bearing bore structure would lead to decrease of the amplitude of the power turbine disc.

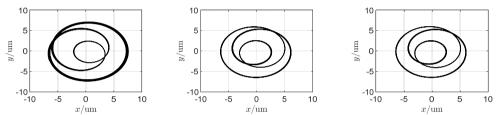


Fig. 4. Influence of radial stiffness of the shared bearing bore-orbit of the power turbine

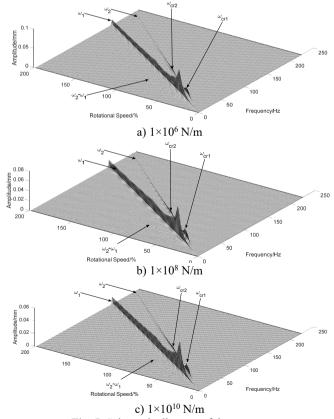


Fig. 5. Schematic diagram of the support

3.2. Influence of the angular stiffness of the shared bearing bore structure

Similar to Section 3.1, different values of the angular stiffness of the shared bearing bore structure were taken with the other parameters remain unchanged. The orbit of the turbine disc at 5400 r/min is shown in Fig. 6 and the waterfall diagram is shown in Fig. 7.

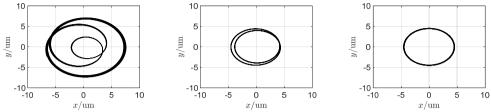


Fig. 6. Influence of angular stiffness of the shared bearing bore-orbit of the power turbine

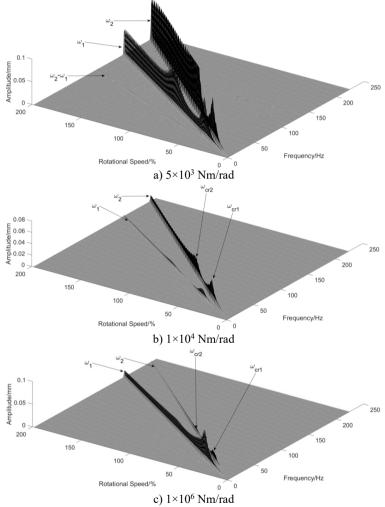


Fig. 7. Influence of angular stiffness of the shared bearing bore

It can be seen from Fig. 7 that the "cross excitation" phenomenon was gradually weakened with increasing angular stiffness of the shared bearing bore. This was also indicated by the orbit

plot shown in Fig. 6. The orbit of the power turbine disc at 5400 r/min gradually forms only one circle, Fig. 6(c), instead of multiple circles, Fig. 6(a). Thus, the angular stiffness of the shared bearing bore structure can effectively reduce the cross-excitation phenomenon as well as the vibration transmission between both rotors.

4. Conclusions

The dynamic characteristics of a coaxial-rotor system with the shared bearing bore structure was studied. The influences of the radial and angular stiffness of the shared bearing bore were simulated and analyzed. The main conclusions are as follows:

- (1) The cross excitation phenomenon was observed in coaxial rotor system with shared bearing bore structure for the first time, indicating that the shared bearing bore structure plays a similar role like the inter-shaft bearing in commercial turbofan aeroengine. The vibration of one rotor would be transmitted to another one by the shared bearing bore thus to increase the opportunity of resonance of the rotor system.
- (2) By numerical simulation it was found that both the radial and angular stiffness of the shared bearing bore has influences on the dynamic response of the rotor system. However, increasing the radial stiffness has little effect on weakening the cross excitation phenomenon. In contrast, increasing the angular stiffness would greatly reduce the cross excitation frequency component, thus the vibration amplitude.

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Data availability

The datasets generated during and/or analyzed during the current study are available from the corresponding author on reasonable request.

Conflict of interest

The authors declare that they have no conflict of interest.

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