On the issue of impact of anisotropy of the rotor elastic suspension on the performance of the automatic balancer

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Abstract. Experimental researches of a vertical rotor with a ball-type automatic balancer with a circular cross-section of the body have proved its efficiency in the reduction of vibrations of an out-of-balance rotor, however for acceleration of compensating masses till the rotor working velocity it is necessary to set initial velocity. A value of the initial velocity of compensating masses depends on the stiffness of the elastic rotor suspension. At the same time, real elastic suspensions of rotors possess nonhomogeneous elastic dissipative characteristics in directions. The research is aimed at determining the influence of anisotropy of the elastic rotor suspension on the acceleration of compensating masses. The authors give recommendations on the selection of parameters of a rotor suspension.

Keywords: automatic balancer, working regime, ball, rotor, torus-shaped body, elastic suspension stiffness.

1. Introduction

A ball-type autobalancing device (hereinafter ABD) allows reducing the disbalance of a rotating out-of-balance rotor on an elastic suspension in the superresonance zone of frequencies. In terms of sensitivity the most efficient is the ABD with a torus-shaped body with the circular cross-section, where circular-shaped compensating masses (hereinafter CM) are freely displaced. However, the acceleration of compensating masses till velocity of the rotor's rotation depends on many parameters of the rotor systems, among which we can mention stiffness of the elastic suspension of the rotor, a coefficient of friction of CM rolling on ABD inner surface [1].

At the same time, for the acceleration of ABD compensating masses in certain cases it is necessary to provide initial velocity [2], otherwise CM steadily move towards ABD body (unsteady operating regime), which leads to increased rotor vibrations. It is obvious that the damping coefficient depends on the suspension stiffness; however, the impact of parameters of the elastic suspension on CM acceleration has been studied insufficiently. Besides, there is no information on the influence of anisotropic elastic suspension of a rotor on the acceleration of ABD compensating masses. In the work [3, 6] we studied the constructions and stability of the autobalancing regime depending on parameters of viscous friction in ABD and in the suspension of the rotor system with a horizontal rotation axis. Therefore, there is certain interest in determining the influence of parameters of the elastic dissipative rotor suspension on the performance of the torus-shaped autobalancing device with a vertical rotation axis in the regime of CM acceleration. The results of the research can be used at the design of an elastic suspension of rotor systems with ABD having torus-shaped body.

The aim of the research is to analyze the process of acceleration of compensating masses of ABD fixed on the rotor with vertical rotation axis on the anisotropic elastic suspension and to determine the conditions for its reliable acceleration. This work is the continuation of work [1] and is based on experimental trials of rotors with ABD [4].

2. Calculation scheme and mathematical model of a rotor with an automatic balancer

The authors study a symmetric vertical stiff rotor on the anisotropic elastic dissipative suspension and with ABD. The body of the automatic balancer is torus-shaped with radius R on the average circumference and R_1 in the cross-section where one ball-shaped compensating mass with radius r and mass m is placed. The rotor with mass m_r possesses static eccentricity e. The rotor is fixed on the anisotropic elastic suspension, whose coefficients of stiffness and viscous damping in the directions of axes x and y equal c_1 , b_1 and c_2 , b_2 accordingly. During the research it was accepted that the value of dissipation coefficients b_1 and b_2 is proportional to the stiffness of the suspension in the corresponding direction. The body of the ABD is concentrically fixed on the rotor and rotates together with it with angular velocity ω .

The calculation scheme of the rotor with ABD shown in Fig. 1 is described in detail in [1]. The system of differential equations of movement describing the movement of the rotor system with ABD, which are given in [1], are specified according to the task.



Fig. 1. Calculation scheme for the rotor system with ABD

Taking into account that rotation velocity is constant ($\dot{\phi} = \omega = \text{const}, \phi = \omega t$), and for the convenience of numerical solution of the set of equations [1] we will present parameters of the rotor system in a dimensionless aspect:

$$\frac{\omega t}{2\pi} = \frac{\varphi}{2\pi} = \tau, \qquad \frac{b_1}{M\omega} = n_1, \quad \frac{c_1}{M\omega^2} = p_1, \quad \frac{b_2}{M\omega} = n_2, \quad \frac{c_2}{M\omega^2} = p_2, \quad \frac{m}{M} = m_0, \\ \frac{M-m}{M} = m_r = 1 - m_0, \quad \frac{R_1 - r}{e} = e_0, \quad \frac{R_1}{e} = e_1, \quad \frac{R}{e} = e_2, \quad \frac{k}{r} = k_r, \quad g_0 = \frac{9,81}{e\omega^2}, \\ b_1 = \mu \cdot c_1, \quad b_2 = \mu \cdot c_2, \end{cases}$$

where τ – dimensionless time.

The task is solved at the constant value of the coefficient of viscous damping (b = const) and with the coefficient proportional to the stiffness of the suspension (μ – proportionality coefficient).

At the numerical solution of differential equations we accept e = 1 m and $\omega = 1$ s⁻¹. At doing the practical task for obtaining real values of the calculation results, they are to be multiplied by specific values e and ω :

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$$\begin{aligned} \frac{\ddot{x}}{e\omega^2} &= -n_1 \frac{\dot{x}}{e\omega} - p_1 \frac{\dot{x}}{e} + (1 - m_0) \cos 2\pi \tau + m_0 e_2 \left(1 + \frac{\dot{\psi}}{\omega}\right)^2 \cos(2\pi \tau + \psi) \\ &+ m_0 e_2 \frac{\ddot{\psi}}{\omega^2} \sin(2\pi \tau + \psi) - m_0 e_0 \frac{\ddot{a}}{\omega^2} \cos(2\pi \tau + \psi) \cos \alpha \end{aligned} \tag{1} \\ &+ m_0 e_0 \left(\frac{\dot{a}}{\omega}\right)^2 \cos(2\pi \tau + \psi) \sin \alpha + 2m_0 e_0 \frac{\dot{a}}{\omega} \left(1 + \frac{\dot{\psi}}{\omega}\right) \sin(2\pi \tau + \psi) \cos \alpha, \end{aligned} \tag{1} \\ &- m_0 e_0 \left(\frac{\dot{a}}{\omega}\right)^2 \cos(2\pi \tau + \psi) - m_0 e_0 \frac{\ddot{a}}{\omega^2} \sin(2\pi \tau + \psi) \cos \alpha \end{aligned} \tag{2} \\ &+ m_0 e_0 \left(\frac{\dot{a}}{\omega}\right)^2 \sin(2\pi \tau + \psi) \sin \alpha - 2m_0 e_0 \frac{\dot{a}}{\omega} \left(1 + \frac{\dot{\psi}}{\omega}\right) \cos(2\pi \tau + \psi) \cos \alpha, \end{aligned} \tag{2} \\ &+ m_0 e_0 \left(\frac{\dot{a}}{\omega}\right)^2 \sin(2\pi \tau + \psi) \sin \alpha - 2m_0 e_0 \frac{\dot{a}}{\omega} \left(1 + \frac{\dot{\psi}}{\omega}\right) \cos(2\pi \tau + \psi) \cos \alpha, \end{aligned} \tag{2} \\ &+ m_0 e_0 \left(\frac{\dot{a}}{\omega}\right)^2 \sin(2\pi \tau + \psi) \sin \alpha - 2m_0 e_0 \frac{\dot{a}}{\omega} \left(1 + \frac{\dot{\psi}}{\omega}\right) \cos(2\pi \tau + \psi) \cos \alpha, \end{aligned} \tag{3} \\ &- \frac{1}{e_0} \left[\frac{1 + \frac{\dot{\psi}}{\omega}}{e} \cos(2\pi \tau + \psi) - \frac{n_1 \dot{x}}{e} \sin(2\pi \tau + \psi)\right] \\ &- \frac{1}{e_0} \left[\frac{1 + \frac{\dot{\psi}}{\omega}}{e}\right] \frac{\dot{a}}{\omega} (1 - m_0) \cos \alpha + \frac{(1 - m_0) \sin \psi}{e_2} - \frac{N^* k_r (e_2 + e_1 \sin \alpha)}{e_2^2} \sin(2\pi \tau + \psi)\right] \\ &- \frac{\ddot{a}}{\omega^2} = \frac{1}{\left(\frac{7}{5} - m_0 \cos^2 \alpha\right)} \left\{\frac{\cos \alpha}{e_0} \left[\frac{n_1 \dot{x}}{e\omega} \cos(2\pi \tau + \psi) + \frac{n_2 \dot{y}}{e\omega} \sin(2\pi \tau + \psi)\right] \\ &- m_0 \left(\frac{\dot{a}}{\omega}\right)^2 \sin \alpha \cos \alpha + \frac{\cos \alpha}{e_0} \left[p_1 x \cos(2\pi \tau + \psi) + p_2 y \sin(2\pi \tau + \psi)\right] - \frac{g_0 \sin \alpha}{e_0} \end{aligned} \tag{4} \\ &- \frac{(1 - m_0) \cos \alpha \cos \psi}{e_0} + \frac{e_2 \cos \alpha}{e_0} \left(1 + \frac{\dot{\psi}}{\omega}\right)^2 (1 - m_0) - \frac{N^* k_r e_1}{e_0^2} \sin \alpha^2 \right\}, \end{aligned}$$

where N^* – dimensionless force of normal pressure of the ball onto the inner surface of ABD body:

$$N^* = g_0 \cos\alpha + e_2 \left(1 + \frac{\dot{\psi}}{\omega}\right)^2 \sin\alpha + e_0 \dot{\alpha}^2.$$
(5)

The calculation results of the system of differential Eqs. (1-4) are obtained in the form of dimensionless velocities and coordinates:

 $v_1 = \frac{\dot{x}}{e\omega}, \ x_1 = \frac{x}{e}, \ v_2 = \frac{\dot{y}}{e\omega}, \ x_2 = \frac{y}{e}, \ v_3 = \frac{\dot{\psi}}{\omega}, \ x_3 = \psi, \ v_4 = \frac{\dot{\alpha}}{\omega}, \ x_4 = \alpha.$

3. Research results of the mathematical model

The differential Eqs. (1-4) of movement of the rotor system with ABD were investigated by means of SPRING software [5].

The authors research the influence of asymmetric stiffness of the elastic rotor suspension on the conditions of CM acceleration till velocity of rotor rotation. Thus, we determined the initial velocity of CM in a circumferential direction of the balancer $v_3 = \dot{\psi}/\omega$ (at the initial conditions CM $\tau = 0$, $x_3 = 0$, $v_4 = 0$, $x_4 = 0$) at which its acceleration till working velocity of the rotor is provided. The initial conditions of the rotor at the study of CM acceleration are accepted as

follows: $\tau = 0, x_1 = 0, v_1 = 0, x_2 = 0, v_2 = 0, \omega = 1.$

At the first stage of the research it was accepted that the dissipation coefficient in the elastic rotor suspension does not depend on stiffness of the suspension and remains constant. The trials of CM acceleration were performed at two fixed values of the coefficient of relative dissipation: $n_1 = n_2 = 0,35$ and $n_1 = n_2 = 0,035$ and at the following parameters of the rotor systems: $m_0 = 0,005$; $e_0 = 2,5$; $e_1 = 15$; $e_2 = 200$; $k_r = 0,0024$; $g_0 = 0,5005$. The calculation results are given in Table 1 and were used for construction of charts in Fig. 2 showing the dependence of the initial velocity of CM v_3 in the circumferential direction of ABD body on the value of dimensionless stiffness of anisotropic rotor suspension. The trials were carried out at a ratio of working rotor velocity to fundamental frequencies of the elastic suspension of the rotor system within: $\omega/\omega_{1,2} = 3,194...22,36$ and $\omega_2/\omega_1 = 1,...,7$.

n = 0.035		The suspension dimensionless stiffness, p_1					
<i>n</i> = 0,035		0,0020	0,0308	0,0500	0,0692	0,0980	
	0,0020	-0,9075	-0,7638	-0,7083	-0,6626	-0,6043	
The suspension	0,0308	-0,7600	-0,7359	-0,6970	-0,6562	-0,6006	
dimensionless stiffness, p_2	0,0500	-0,7035	-0,6945	-0,6766	-0,6466	-0,5962	
	0,0692	-0,6571	-0,6522	-0,6444	-0,6281	-0,5884	
	0,0980	-0,5979	-0,5952	-0,5918	-0,5857	-0,5663	
m = 0.25		The suspension dimensionless stiffness, p_1					
m = 0.35		The	suspension	dimension	less stiffnes	s, <i>p</i> ₁	
<i>n</i> = 0,35		The 0,0020	suspension 0,0308	dimension 0,0500	less stiffnes 0,0692	s, p ₁ 0,0980	
n = 0,35	0,0020	The 0,0020 -1,000	suspension 0,0308 0,4888	dimension 0,0500 -0,4334	less stiffnes 0,0692 -0,3884	s, p ₁ 0,0980 -0,3313	
n = 0,35 The suspension	0,0020 0,0308	The 0,0020 -1,000 -0,4887	suspension 0,0308 -0,4888 -0,4199	dimension 0,0500 -0,4334 -0,3813	ess stiffnes 0,0692 -0,3884 -0,3463	s, p ₁ 0,0980 -0,3313 -0,2984	
n = 0,35 The suspension dimensionless stiffness,	0,0020 0,0308 0,0500	The 0,0020 -1,000 -0,4887 -0,4332	suspension 0,0308 -0,4888 -0,4199 -0,3812	dimension 0,0500 -0,4334 -0,3813 -0,3493	ess stiffnes 0,0692 -0,3884 -0,3463 -0,3191	s, p ₁ 0,0980 -0,3313 -0,2984 -0,2763	
n = 0.35 The suspension dimensionless stiffness, p_2	0,0020 0,0308 0,0500 0,0692	The 0,0020 -1,000 -0,4887 -0,4332 -0,3882	suspension 0,0308 -0,4888 -0,4199 -0,3812 -0,3461	dimension 0,0500 -0,4334 -0,3813 -0,3493 -0,3190	ess stiffnes 0,0692 -0,3884 -0,3463 -0,3191 -0,2926	s, p ₁ 0,0980 -0,3313 -0,2984 -0,2763 -0,2541	

Table 1. The initial velocity of the compensating mass in the circumferential direction of the balancer v_3





Table 1 and the charts (Fig. 2) make it clear that anisotropy of the elastic rotor suspension influences the value of CM initial velocity. The increase of the suspension stiffness (parameter "p") in any direction leads to the necessity to increase absolute initial velocity of the compensating mass for its acceleration till the working velocity of the rotor. Vice versa, the reduction of stiffness of the rotor suspension in any direction allows decreasing the absolute velocity of CM. Thus, the increase of the dissipation coefficient in the rotor suspension leads to the necessity to increase the

absolute velocity of CM for its acceleration till rotor velocity.

At the second stage of the research it was accepted that the dissipation coefficient in the elastic rotor suspension is proportional to the stiffness of the suspension, i.e. $n_1 = \mu \cdot c_1$; $n_2 = \mu \cdot c_2$. The trials of CM acceleration were carried out at two values of proportionality coefficient ($\mu = 0.7$ and $\mu = 7.0$), which correspond to the values of relative dissipation in Table 1 at p = 0.05 (average value in a range of relative stiffness). Other parameters of the rotor system did not change. The research results are summarized in Table 2 and in Fig. 3. As may be seen from the analysis of Table 2 and Fig. 3 anisotropy of the dissipation coefficient of the suspension makes essential impact on CM acceleration in ABD.

u = 0.7		The suspension dimensionless stiffness, p_1				
$\mu = 0, 7$		0,0020	0,0308	0,0500	0,0692	0,0980
	0,0020	-0,9355	-0,7774	-0,7129	-0,6586	-0,5881
The suspension	0,0308	-0,7714	-0,7519	-0,7018	-0,6506	-0,5815
dimensionless stiffness,	0,0500	-0,7078	-0,6992	-0,6766	-0,6370	-0,5729
p_2	0,0692	-0,6540	-0,6475	-0,6356	-0,6115	-0,5591
	0,0980	-0,5841	-0,5783	-0,5708	-0,5578	-0,5249
		The suspension dimensionless stiffness, p_1			s, p_1	
$\mu = 7,0$		0,0020	0,0308	0,0500	0,0692	0,0980
	0,0020	-0,9420	-1,0000	-1,0000	-1,0000	-1,0000
The suspension	0,0308	-0,9873	-0,5464	-0,4633	-0,4191	-0,4422
dimensionless stiffness,	0,0500	-0,9874	-0,4640	-0,3493	-0,2696	-0,2054
p_2	0,0692	-0,9872	-0,4204	-0,2701	-0,1612	-0,0395
	0,0980	-0,9865	-0,4442	-0,2064	-0,0400	-0,0000
-0,55	0,098	s, p2				0,0980

	Table	2. The initial	velocity of the	compensating ma	ass in the circur	nferential directio	n of the balancer v_3
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4. Conclusions

In the result of performed calculations of the mathematical model for a vertical rotor with concentrically fixed torus-shaped ABD body with a circular cross-section, it has been stated that CM acceleration till rotor working velocity essentially depends on anisotropy of stiffness and dissipation coefficients of the elastic rotor suspension. Thus, the increase of anisotropy of the elastic rotor suspension towards the reduction of stiffness coefficient in one direction and the increase of dissipation leads to the reduction of necessary initial absolute CM velocity in the circumferential direction ABD for its acceleration till rotor velocity. For successful CM acceleration it is recommended to introduce temporary (for the time of acceleration) artificial anisotropy into the design of a rotor suspension.

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