

A weak solution for free vibration of multi-span beams with general elastic boundary and coupling conditions

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Abstract. A weak solution of free vibration is developed for multi-span beams, which can adapt general elastic boundary and coupling conditions. Firstly, create the energy functional of the multi-span beam system based on the small deformation theory. Then, adopt the modified Fourier series method to rewrite the displacement functions. Compared with the traditional Fourier series method, the present series representations provide a solution for general elastic restraints. Lastly, combined with the Rayleigh-Ritz technique, all the series expansion coefficients can be obtained as the generalized coordinates. Numerical results demonstrate that the current weak solution has good convergence and high accuracy compared with the existing results in literature and FEM results.

Keywords: weak solution, free vibration, multi-span beams, general boundary conditions, elastic coupling conditions.

1. Introduction

Beams have been applied in many filed such as aerospace, marine, construction industry due to their excellent engineering features. And in some specific field, the multi-span beam structures are also used widely, such as bridge engineering, water conservancy projects, etc. As a result, over the years, the vibration of the beam and multi span beam structure characteristics has always been the hot spot of the scholars' concerned problem.

As so far, many scholars have proposed a number of computational techniques for multi-span beams. Zheng et al. [1] apply the assumed modes method to study the vibration of a multi-span non-uniform beam subjected to a moving load. Wang [2] presents a method of modal analysis to investigate the forced vibration of multi-span Timoshenko beams with clamped boundary condition and rigid joint. The vibration of a multi-span non-uniform bridge subjected to a moving vehicle is analyzed by Cheung et al. [3] using the assumed modes method. Based on the modal analysis method and the direct integration method, Dugush et al. [4] and Ichikawa [5] study the dynamic behavior of multi-span non-uniform beams. Lin and Chang [6] deal with the free vibration analysis of a multi-span beam with an arbitrary number of flexible constraints by using the transfer matrix method. Li and Xu [7] present an exact Fourier Series method to study the vibration analysis of multi-span beam systems. Lin and Tsai [8] use the finite element method to analyze free vibration analysis of a uniform multi-span beam carrying multiple spring-mass systems. Hong and Kim [9] present an exact modelling and modal analysis method for non-uniform, multi-span beam-type structure supported and/or connected by resilient joints with damping. Marchesiello [10] presents an analytical approach to study the dynamics of multi-span continuous straight bridges subject to multi-degrees of freedom moving vehicle excitation. Johansson et al. [11] present a closed-form solution for evaluating the dynamical behavior of a general multi-span Bernoulli–Euler beam by using mode superposition method. Lin [12] employs the numerical assembly method (NAM) to determine dynamic behavior of a multi-span uniform beam carrying a number of various concentrated elements.

In this paper, the authors present a weak solution for free vibration of multi-span beams

subjected to general elastic boundary and coupling conditions. The general elastic boundary and coupling constraints of the multi-span beams are realized by applying the artificial stiffness-like spring technique. Based on the modified Fourier series method, each of displacements is written as a form of a standard one-dimensional Fourier cosine series with several auxiliary functions. The series expansion coefficients are obtained by using the Rayleigh-Ritz technique. The accuracy and convergence of the current weak solution is checked by comparing with the existing results in literature and FEM results.

2. Theoretical formulations

A multi-span beam system is shown as Figure 1. The system includes multiple beams which are coupled with joints. Each joint consists of two group springs which are linear and rotational springs respectively. The effects of non-rigid or resilient connectors are allowed in using the coupling springs between two adjacent beams. When the stiffness of these springs become considerably larger than the bending rigidities of the involved beams, it turns to be a conventional rigid connector. Each of beams are supported by a set of elastic restraints at both ends. By adjusting the spring stiffness values from zero to infinite, different boundary conditions including traditional intermediate supports and classical boundary conditions (i.e., the combinations of the simply supported (S), free (F), and clamped end conditions (C)) can be obtained.

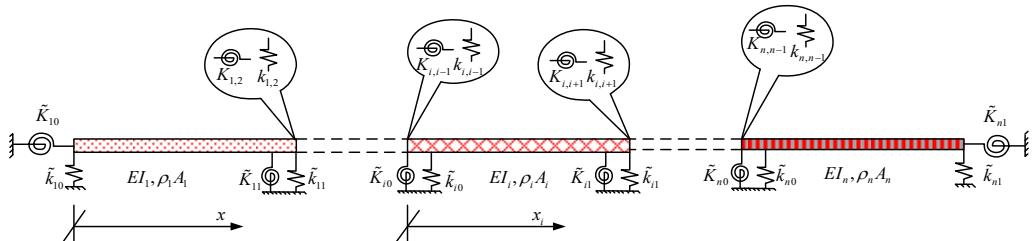


Fig. 1. Multi-span beam system with general boundary conditions

The differential equation for the vibration of the i th beam can be expressed as:

$$D_i d^4 w_i(x) / dx^4 - \rho_i A_i \omega^2 w_i(x) = 0, \quad (1)$$

where w is the angular frequency of the beam, w_i , D_i , A_i and ρ_i are the displacement function, bending rigid, the cross-sectional area and the density of the beams, respectively.

From the previous reviews, in this study, the artificial stiffness-like spring technique is adopted to simulate the arbitrary boundary conditions and continuity conditions. With this method, the boundary and continuity conditions can be expressed as follows:

At $x_i = 0$:

$$k_{i,i-1}(w_i - w_{i-1}) + \tilde{k}_{i0}w_i = -D_i w_i''', \quad (2)$$

$$K_{i,i-1}(w'_i - w'_{i-1}) + \tilde{K}_{i0}w'_i = D_i w_i''. \quad (3)$$

At $x_i = L_i$:

$$k_{i,i+1}(w_i - w_{i+1}) + \tilde{k}_{i1}w_i = D_i w_i''', \quad (4)$$

$$K_{i,i+1}(w'_i - w'_{i+1}) + \tilde{K}_{i1}w'_i = -D_i w_i''. \quad (5)$$

At the left end (of the first beam):

$$\tilde{k}_{10}w_1 = -D_1 w_1''', \quad (6)$$

$$\tilde{K}_{10}w'_1 = D_1w''_1. \quad (7)$$

At the right end (of the N th beam):

$$\tilde{k}_{n1}w_n = D_nw'''_n, \quad (8)$$

$$\tilde{K}_{n1}w'_n = -D_nw''_n, \quad (9)$$

where refer to Fig. 1, $k_{i,j}$ denote the stiffnesses of the linear coupling springs in z_i -directions, and $K_{i,j}$ denote the stiffnesses of the rotational coupling springs at the junction of beams i and j , respectively; \tilde{k}_{i0} and \tilde{k}_{i1} are the stiffnesses of linear boundary springs, and \tilde{K}_{i0} , \tilde{K}_{i1} the stiffnesses of the rotational boundary springs at the left and right ends of beam i , respectively.

For the multi-span beams, the total strain energy (V) and kinetic energy (T) can be expressed as:

$$V = \sum_{i=1}^N V_{b,i} + \sum_{i=1}^{N-1} V_{i,i+1}^s, \quad (10)$$

$$T = \sum_{i=1}^N T_{b,i}, \quad (11)$$

where $V_{b,i}$ and $T_{b,i}$ are represent the strain energy and kinetic energy of the i th beams, and $V_{i,i+1}^s$ is the potential energy expression in the connective springs related to i th and $i+1$ th beams. The detailed expression of the $V_{b,i}$, $V_{i,i+1}^s$ and $T_{b,i}$ can be written as:

$$V_{b,i} = \frac{1}{2} \int_0^{L_i} \left\{ \int_0^{L_i} D_i (\partial^2 w_i(x) / \partial x^2) dx \right\} d\theta_i + \frac{1}{2} \left((\tilde{k}_{i0} w_i(x)^2 + \tilde{K}_{i0} (\partial w_i(x) / \partial x)^2)_{x_i=0} + (\tilde{k}_{i1} w_i(x)^2 + \tilde{K}_{i1} (\partial w_i(x) / \partial x)^2)_{x_i=L_i} \right), \quad (12)$$

$$V_{i,i+1}^s = \frac{1}{2} \left\{ +K_{i,i-1} \left(\partial w_i(x) / \partial x \Big|_{x_i=0} - \partial w_{i-1}(x) / \partial x \Big|_{x_{i-1}=L_{i-1}} \right)^2 + K_{i,i+1} \left(w_i(x) \Big|_{x_i=L_i} - w_{i+1}(x) \Big|_{x_{i+1}=0} \right)^2 + K_{i,i+1} \left(\partial w_i(x) / \partial x \Big|_{x_i=L_i} - \partial w_{i+1}(x) / \partial x \Big|_{x_{i+1}=0} \right)^2 \right\}, \quad (13)$$

$$T_{b,i} = \frac{1}{2} \rho_i A_i \omega^2 \int_0^{\phi_i} \{w_i(x)\} R_i d\theta_i. \quad (14)$$

The Lagrangian for the multi-span beams can be generally expressed as:

$$L = T - V. \quad (15)$$

The traditional Fourier series is a well-known form of admissible function for its excellent convergence. However, it is only available to some very simple boundary conditions and would a modified Fourier series technique proposed by Li [13, 14] is widely used in the vibrations of plates and shells with different boundary conditions by Rayleigh-Ritz method. Thus, in this formulation, the modified Fourier series technique is adopted and extended to investigate the free vibrations of multi-span beams with general elastic boundary and coupling conditions:

$$w_i(x_i) = \sum_{m=0}^{\infty} A_{i,m} \cos \lambda_{i,m} x + \sum_{n=1}^4 B_{i,n} \sin \lambda_{i,n} x, \quad (\lambda_m = m\pi/L_i), \quad (16)$$

where $\lambda_m = m\pi/L_i$, $A_{i,m}$ and $B_{i,m}$ are the unknown Fourier coefficients of one-dimensional Fourier series expansions for the displacements functions, respectively.

Then, substituting Eq. (16) into the Lagrangian function Eq. (15) and taking its derivatives with respect to each of the undetermined coefficients and making them equal to zero:

$$\frac{\partial L}{\partial \Xi} = 0, \quad \begin{cases} \Xi = A_{i,m}, B_{i,m}, \\ i = 1, 2, \dots, N, \\ m = 1, 2, \dots, M. \end{cases} \quad (17)$$

The problem will be transformed into an eigenvalue and eigenvector problem, and the following governing eigenvalue equation in the matrix form can be achieved:

$$(\mathbf{K} - \omega^2 \mathbf{M}) \mathbf{G} = 0, \quad (18)$$

where \mathbf{K} is the stiffness matrix for the beam, and \mathbf{M} is the mass matrix.

For conciseness, the detailed expression for sub-stiffness and sub-mass matrices will not be shown here. By solving the Eq. (18), the frequencies (or eigenvalues) of multi-span beams can be readily obtained and the mode shapes can be yielded by substituting the corresponding eigenvectors into series representations of displacement components.

3. Case studies

Free vibration of three-span beams with elastic boundary and coupling conditions are researched through the present weak solutions. The related geometrical and material parameters of the following case are given in Table 1.

Table 1. A list of beam parameters and material properties

Variables	Beam 1	Beam 2	Beam 3
L (m)	1.0	1.5	2.0
A (m^2)	5×10^{-5}	1.5×10^{-5}	5×10^{-5}
I (m^4)	10^{-10}	5×10^{-11}	10^{-10}
E (GPa)	207	207	207
ρ (kg/m^3)	7800	7800	7800

The schematic diagram, geometrical and material parameters of three-span beam are shown in Fig. 2 and Table 1, respectively. Table 2 gives the corresponding stiffness of the boundary and coupling springs with refer to the Fig. 2. Also, the Table 3 gives the convergence and validation study. In the FEM mode, the element type and mesh size are the beam 188 and 0.1 m, respectively. The mode shapes for the first four modes are plotted in Fig. 3. From the Table 3 and Fig. 3, it is obvious that the present weak solution is not only available to solve the multi-span beams with elastic boundary and coupling conditions, but also has a good accuracy and reliability.

Table 2. Stiffness values for the boundary and coupling springs of a three-beam system

Spring constants for beam 1	Spring constants for beam 2	Spring constants for beam 3	Spring constants for joints
$\tilde{k}_{10} = 10^{10}$ N/m $\tilde{k}_{11} = 5000$ N/m	$\tilde{k}_{20} = 4000$ N/m $\tilde{k}_{21} = 4000$ N/m	$\tilde{k}_{30} = 5000$ N/m $\tilde{k}_{31} = 10^{10}$ N/m	$k_{1,2} = 1000$ N/m $k_{2,3} = 1000$ N/m
$\tilde{K}_{10} = 10^{10}$ Nm/rad	$\tilde{K}_{20} = 1000$ Nm/rad	$\tilde{K}_{30} = 2000$ Nm/rad	$K_{1,2} = 200$ Nm/rad
$\tilde{K}_{10} = 2000$ Nm/rad	$\tilde{K}_{21} = 1000$ Nm/rad	$\tilde{K}_{31} = 0$ Nm/rad	$K_{2,3} = 200$ Nm/rad

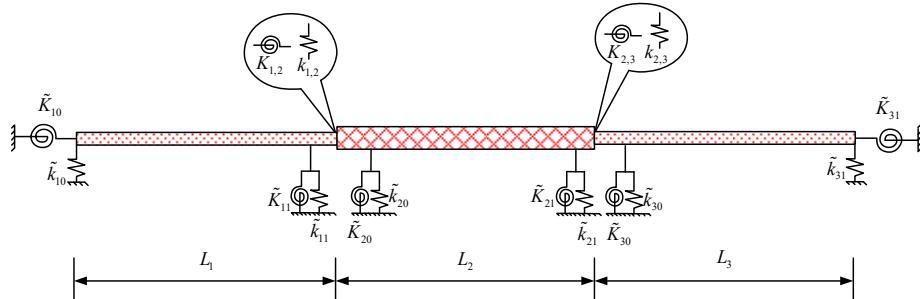


Fig. 2. Schematic diagram of three-span beam

Table 3. The six lowest natural frequencies for various numbers of terms in Fourier series under the three-span beam

Mode	Present Method					Exact solution	FEM
	$M = 6$	$M = 8$	$M = 10$	$M = 12$	$M = 15$		
1	4.3676	4.3675	4.3675	4.3675	4.3675	4.3675	4.3675
2	13.640	13.640	13.639	13.639	13.639	13.646	13.646
3	13.833	13.833	13.833	13.833	13.833	13.848	13.848
4	21.690	21.690	21.690	21.690	21.688	21.689	21.689
5	26.698	26.696	26.696	26.696	26.695	26.697	26.697
6	34.327	34.323	34.323	34.322	34.322	34.325	34.325

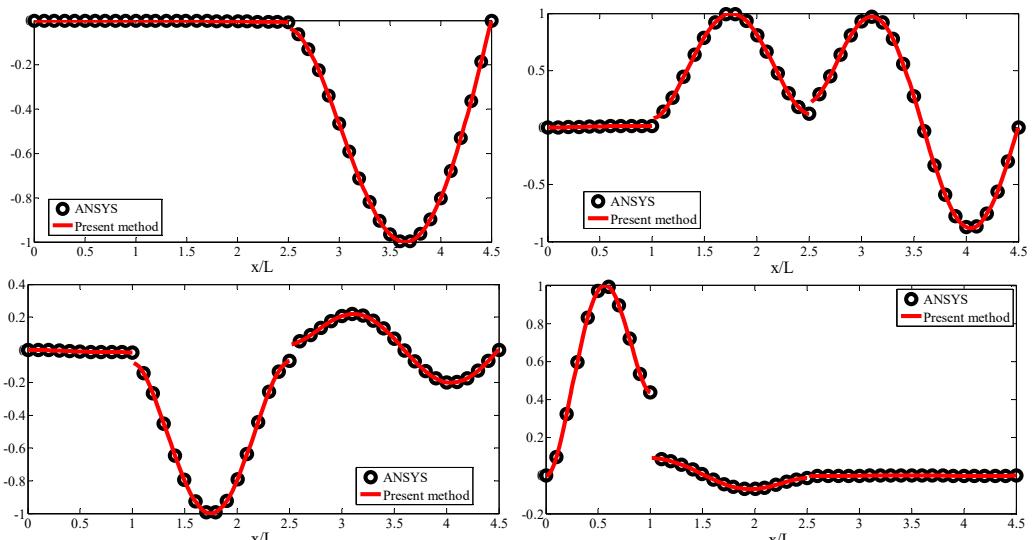


Fig. 3. The four-mode shape for the three-span beam

4. Conclusions

In this paper, a weak solution of free vibration is developed for multi-span beams, which can be applied with general elastic boundary and coupling conditions. Each of displacements is written as a form of a standard one-dimensional Fourier cosine series with several auxiliary functions. The introducing of these functions is to eliminate the discontinuities of all the related displacements and their derivatives at the ends and to improve the convergence speed of the series representations. Combined with the Rayleigh-Ritz technique, all the series expansion coefficients can be obtained as the generalized coordinates. The numerical results show that the current weak solution has good convergence and high accuracy compared with existing results in literature and

FEM results.

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