

# 2404. The vibration frequency analysis of the screw feeding system with spring supports

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Received 23 May 2016; received in revised form 9 November 2016; accepted 24 November 2016  
DOI <https://doi.org/10.21595/jve.2016.17198>



**Abstract.** Ball screw feeding system has wide application in machine tools, and its dynamic research is very important. The drive screw can be seen as a rotating Timoshenko beam under the pre-tension force. Considering the bearing stiffness, the boundary condition of the screw is spring supported. Also, considering the contact deformation between the screw and the worktable and the worktable moving, the screw is added an intermediate moving spring-mass system. Using boundary conditions and continuity conditions, the lateral and longitudinal frequency equations of the screw feeding system are derived considering the effect of gyroscope, the pre-tension force, the bearing stiffness, the contact between the worktable and the screw and the moving of the worktable. The effect of the system parameters on the natural frequency is studied. The diameter, the pre-tension force, the bearing stiffness, the contact stiffness and the rotational speed of the screw have effect on the vibration frequency and the frequency varies when the worktable moves along the screw from the left to the right end. This work supplies a base for designing the drive system to avoid resonance and to improve the stability of the system.

**Keywords:** screw feeding system, vibration, frequency, spring supports.

## Nomenclature

$t$	Time
$v(x, t)$	Lateral displacement of the screw in $y$ direction
$\psi(x, t)$	Rotational angle of the screw around $z$ axis
$u(x, t)$	Axial displacement of the screw in $x$ direction
$s_v$	Displacement of the worktable in $y$ direction
$s_u$	Displacement of the worktable in $x$ direction
$\phi_m$	Rotational angle of the worktable around $z$ axis
$k_{v1}, k_{v2}$	Lateral stiffness coefficients of the screw bearings in $y$ direction
$k_{\psi1}, k_{\psi2}$	Rotational stiffness coefficients of the bearings around $z$ axis
$k_v$	Contact stiffness coefficient between the screw and worktable in $y$ direction
$k_\psi$	Contact stiffness coefficient between the screw and worktable around $z$ axis
$k_{g1}, k_{g2}$	Contact stiffness coefficients between the worktable and the guide way
$k_{u1}, k_{u2}$	Axial stiffness coefficients of the screw bearings in $x$ direction
$k_u$	Contact stiffness coefficient between the screw and nut in $x$ direction
$L$	Screw length
$\lambda$	Position coefficient of the contact point between the worktable and the screw
$E$	Young's modulus of the screw material
$G$	Shear modulus of the screw material
$\rho$	Density of the screw material
$k$	Cross-section shape factor of the screw
$d$	Diameter of the screw
$A$	Cross section area of the screw, $A = \pi d^2/4$
$I$	Second moment of inertia of the screw, $I = \pi d^4/64$
$\Omega$	Rotational speed of the screw

$\omega$	Lateral vibration frequency of the screw feeding system
$\omega_u$	Longitudinal vibration frequency of the screw feeding system
$\omega_{Oss}$	Lateral vibration frequency of Euler-Bernoulli beam under simply supported at both ends
$\omega_{Off}$	Lateral vibration frequency of Euler-Bernoulli beam under freely supported at both ends
$\omega_{Occ}$	Lateral vibration frequency of Euler-Bernoulli beam under clamed supported at both ends
$\omega_{Tss}$	Lateral vibration frequency of Timoshenko beam under simply supported at both ends
$\omega_{Tff}$	Lateral vibration frequency of Timoshenko beam under freely supported at both ends
$\omega_{Tcc}$	Lateral vibration frequency of Timoshenko beam under clamed supported at both ends
$\omega_{uff}$	Longitudinal vibration frequency of the rod under freely supported at both ends
$\omega_{ucc}$	Longitudinal vibration frequency of the rod under clamed supported at both ends
$\omega_{ucf}$	Longitudinal vibration frequency of the rod under freely supported at one end and clamed supported at the other end
$F$	Axial pre-tension force of the screw
$a$	Length of the worktable
$I_z$	Worktable's moment of inertia around z axis, $I_z = ma^2/12$

## 1. Introduction

Ball screw feeding system has wide application in machine tools. With increasing of the feeding velocity and machining precision, the vibration of the feeding system becomes more and more important, and gets attention from researchers.

Yang et al. [1] structured the dynamic model of screw-workbench system based on the Timoshenko beam theory account of lateral shear of beam. Feng et al. [2] analyzed the modal and dynamic response of the feeding screw using Ansys. Li et al. [3] used FEA method to do the modal analysis and studied the un-damped free vibration of the leading screw, and got the intrinsic frequency and the vibration model. Jun [4] analyzed ball screw feeding system's static and dynamic characteristic of a CNC lathe, and got the natural frequency and vibration mode of the feeding system. Dong et al. [5] presented a method to model and design servo controllers for flexible ball screw drives with dynamic variations and proposed a mathematical model describing the structural flexibility of the ball screw drive containing time-varying uncertainties and disturbances with unknown bounds. Feng et al. [6] examined the relationship between the ball screw preload variation and detected vibration signals. Wang et al. [7] took the form of a hollow screw shaft structure containing multiple tuned mass dampers to achieve the lateral multi-mode vibration control of the screw shaft, considering the time-varying vibration characteristics of the ball screw caused by the continuous movement of the nut.

In the screw drive system, the screw is slender, and it can be seen as Timoshenko beam. The worktable moves with the screw rotation, and its mass is much larger than the screw, so the worktable can be seen as a moving mass. Considering contact stiffness between the worktable and the screw, the screw drive system is a rotating beam with moving spring-mass system. Tang et al. [8] derived exact frequency equations of the Timoshenko beam for various boundary conditions. Arboleda-Monsalvea et al. [9] analyzed the stability and free vibration of a Timoshenko beam-column with generalized end conditions subjected to constant axial load, and used intermediate flexural connection to model the crack of the beam. You et al. [10] analyzed natural frequencies, modes and critical speeds of axially moving beams on different supports based on Timoshenko model, and derived the governing differential equation of motion from Newton's second law. Azama et al. [11] analyzed the dynamic responses of a Timoshenko beam subjected

to a moving mass and a moving sprung mass, also derived governing differential equations for beam vibration. Dehestani et al. [12] presented an analytical-numerical method to determine the dynamic response of beams carrying a moving mass, and investigated the influences regarding the speed of the moving mass on the dynamic response of beams. Ghayesh et al. [13] investigated the nonlinear forced vibrations and stability of an axially moving Timoshenko beam with an intra-span spring-support. Lin [14] utilized the numerical assembly method (NAM) to determine the exact natural frequencies and mode shapes of the multi-span Timoshenko beam carrying a number of various concentrated elements including point masses, rotary inertias, linear springs, rotational springs and spring-mass systems. Lee et al. [15] analyzed the dynamic contact problem of a tensioned beam with clamped-pinned ends when the beam contacts a moving mass-spring system, derived the dynamic contact equations using constraints and equations of motion for the beam and moving mass and then discretized using the finite element method. The above research focused on the lateral vibration, but also the longitudinal vibration existed. Cortés et al. [16] studied the longitudinal vibration behaviour of a damped rod.

Zhang et al. [17-20] also studied the vibration and dynamic response of the screw feeding system in consideration of the bearing stiffness, the contact deformation of the feeding system, but neglected the moving of the worktable. In present work, the frequency equation of the screw feeding system under elastically supported is derived, considering the effect of gyroscope, the pre-tension force, the bearing stiffness, the contact stiffness of the system and the intermediate moving spring-mass. Effects of the system parameters on the screw's lateral and longitudinal frequency are discussed.

## 2. Lateral vibration frequency equation of the screw feeding system

Using separation of variables, the lateral vibration displacement of the beam can be written as:

$$\begin{cases} v(x, t) = V(x)e^{i\omega t}, \\ \psi(x, t) = \Psi(x)e^{i\omega t}. \end{cases} \quad (1)$$

From the Ref. [20], considering the shear deformation, rotary inertia and the pretension force, the mode shape of the Timoshenko beam can be expressed as:

$$\begin{cases} V(x) = C_1 \sin(\beta x) + C_2 \cos(\beta x) + C_3 \sinh(\alpha x) + C_4 \cosh(\alpha x), \\ \Psi(x) = \eta C_1 \cos(\beta x) - \eta C_2 \sin(\beta x) + \delta C_3 \cosh(\alpha x) + \delta C_4 \sinh(\alpha x), \end{cases} \quad (2)$$

where:

$$\begin{aligned} \alpha &= \frac{-b_0 + \sqrt{b_0^2 - 4a_0c_0}}{2a_0}, \quad \beta = \frac{b_0 + \sqrt{b_0^2 - 4a_0c_0}}{2a_0}, \\ a_0 &= (GAk + F)EI, \quad b_0 = EI\rho A\omega^2 + (F + GAk)(-GAk + \rho I\omega^2 - 2\rho I\Omega\omega) + (GAk)^2, \\ c_0 &= \rho A\omega^2(\rho I\omega^2 - GAk - 2\rho I\Omega\omega), \\ \eta &= \frac{GAk\beta^2 + F\beta^2 - \rho A\omega^2}{GAk\beta}, \quad \delta = \frac{GAk\alpha^2 + F\alpha^2 + \rho A\omega^2}{GAk\alpha}. \end{aligned}$$

If  $\alpha$  is an imaginary,  $\sin(\alpha)$  and  $\cos(\alpha)$  replace  $\sinh(\alpha)$  and  $\cosh(\alpha)$ , respectively.

The dynamic modal of the screw feeding system is shown in Fig. 1. To consider the contact deformation, the bearings at both ends of the screw are expressed by spring stiffness coefficients. The worktable and the screw connect at point  $O'$ . The contact deformation between the screw and the worktable is also expressed as spring stiffness coefficients, so an intermediate spring-mass system of the screw is added. And also, the contact deformation between the worktable and the guide way is considered using spring stiffness coefficients.

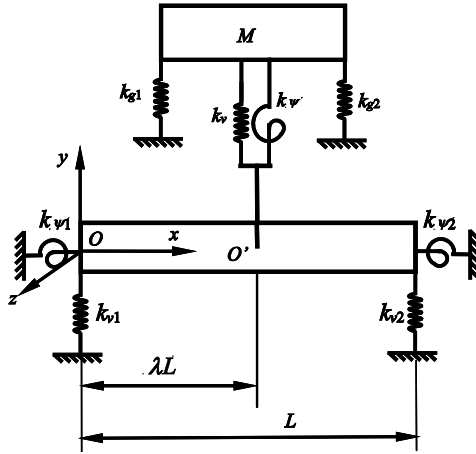


Fig. 1. The dynamic model of the screw feeding system lateral vibration

The screw and the worktable contact at the point  $O'$  and the distance between the contact point and the left end of the screw is  $\lambda L$ . So, the screw is made of two segments. The lateral vibration mode shape of the screw can be written as:

$$\begin{cases} V(x) = \begin{cases} C_1 \sin(\beta x) + C_2 \cos(\beta x) + C_3 \sinh(\alpha x) + C_4 \cosh(\alpha x), & x \in [0, \lambda L], \\ C_5 \sin(\beta x) + C_6 \cos(\beta x) + C_7 \sinh(\alpha x) + C_8 \cosh(\alpha x), & x \in [\lambda L, L], \end{cases} \\ \Psi(x) = \begin{cases} \eta C_1 \cos(\beta x) - \eta C_2 \sin(\beta x) + \delta C_3 \cosh(\alpha x) + \delta C_4 \sinh(\alpha x), & x \in [0, \lambda L], \\ \eta C_5 \cos(\beta x) - \eta C_6 \sin(\beta x) + \delta C_7 \cosh(\alpha x) + \delta C_8 \sinh(\alpha x), & x \in [\lambda L, L], \end{cases} \end{cases} \quad (3)$$

where  $C_1$ - $C_8$  are coefficients of the lateral vibration mode shape yet to be determined.

The bearings at both ends of the screw are expressed by spring's stiffness, so the screw is spring supported at both ends, and the boundary conditions of the lateral vibration are:

$$\begin{cases} EI \frac{\partial \psi(0, t)}{\partial x} = k_{\psi 1} \psi(0, t), \\ kAG \left( \psi(0, t) - \frac{\partial v(0, t)}{\partial x} \right) = -k_{v1} v(0, t), \\ EI \frac{\partial \psi(L, t)}{\partial x} = -k_{\psi 2} \psi(L, t), \\ kAG \left( \psi(L, t) - \frac{\partial v(L, t)}{\partial x} \right) = k_{v2} v(L, t). \end{cases} \quad (4)$$

The screw continues at contact point  $O'$ , and the compatibility conditions at point  $O'$  can be expressed as:

$$\begin{cases} v(\lambda L^-, t) = v(\lambda L^+, t), \\ \psi(\lambda L^-, t) = \psi(\lambda L^+, t), \\ k_v [v(\lambda L, t) - s_v] = kAG \left( \psi(\lambda L^-, t) - \frac{\partial v(\lambda L^-, t)}{\partial x} \right) - kAG \left( \psi(\lambda L^+, t) - \frac{\partial v(\lambda L^+, t)}{\partial x} \right), \\ k_\psi [\psi(\lambda L, t) - \varphi_m] = EI \frac{\partial \psi(\lambda L^-)}{\partial x} - EI \frac{\partial \psi(\lambda L^+)}{\partial x}, \end{cases} \quad (5)$$

where “-” expresses the left side of the contact point  $O'$ , and “+” expresses the right side of the contact point  $O'$ .

The vibration of the worktable includes the vibration in  $y$  direction and rotation around the  $z$  axis, which equations are:

$$M\ddot{s}_v = -k_v(s_v - v) - (k_{g1} + k_{g2})s_v, \quad (6)$$

$$I_z\ddot{\varphi}_m = -(k_{g1} + k_{g2})(\varphi_m - \psi(\lambda L, t))a. \quad (7)$$

Letting:

$$s_v = S_v e^{i\omega t}, \quad (8)$$

$$\varphi_m = \Phi_m e^{i\omega t}. \quad (9)$$

Substituting Eqs. (1), (3), (8) and (9) into Eqs. (4-7), these equations can be arranged as:

$$C[C_1 \ C_2 \ C_3 \ C_4 \ C_5 \ C_6 \ C_7 \ C_8 \ S_v \ \Phi]^T = 0. \quad (10)$$

For the non-trivial solution, the coefficient matrix must be singular, so the frequency equation is:

$$|C| = 0. \quad (11)$$

### 3. The longitudinal vibration frequency equation of the screw feeding system

From Ref. [21], the longitudinal vibration displacement of the rod is written as:

$$u(x, t) = U(x)e^{i\omega_u t}. \quad (12)$$

The longitudinal vibration mode shape of the rod is expressed as:

$$U(x) = D_1 \sin \frac{\omega_u}{a} x + D_2 \cos \frac{\omega_u}{a} x, \quad (13)$$

where  $a^2 = E/\rho$ .

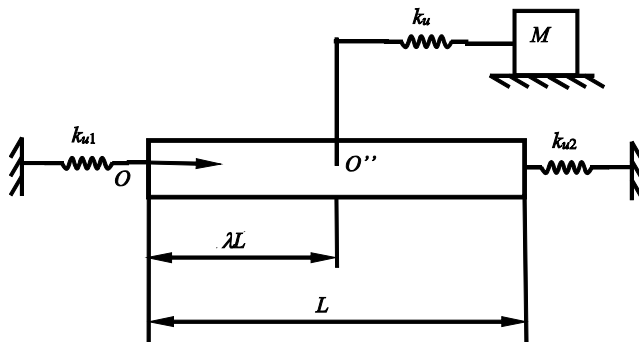


Fig. 2. The dynamic model of the screw feeding system longitudinal vibration

The longitudinal vibration dynamic model of the screw feeding system is shown in Fig. 2. Considering the bearing's stiffness, the screw is spring supported at the both ends. Considering the contact deformation between the ball screw and the moving worktable, the screw is added with an intermediate moving spring-mass system. The distance between the contact point and the left end of the screw is  $\lambda L$ . The longitudinal vibration mode shape of the screw is:

$$\begin{cases} U(x) = D_1 \sin \frac{\omega_u}{a} x + D_2 \cos \frac{\omega_u}{a} x, & x \in [0, \lambda L], \\ U(x) = D_3 \sin \frac{\omega_u}{a} x + D_4 \cos \frac{\omega_u}{a} x, & x \in [\lambda L, L], \end{cases} \quad (14)$$

where  $D_1$ - $D_4$  represent four integration constants yet to be determined.

Both ends of the screw subject to elastic supports, so the boundary conditions of the screw are:

$$\begin{cases} EA \frac{\partial u_1(0, t)}{\partial x} = k_{u1} u(0, t), \\ EA \frac{\partial u_2(L, t)}{\partial x} = -k_{u2} u(L, t). \end{cases} \quad (15)$$

The screw continues at the contact point  $O''$ , and the compatibility conditions are:

$$\begin{cases} u(\lambda L^-, t) = u(\lambda L^+, t), \\ k_u (s_u - u(\lambda L, t)) = EA \frac{\partial u(\lambda L^-, t)}{\partial x} - EA \frac{\partial u(\lambda L^+, t)}{\partial x}. \end{cases} \quad (16)$$

According to the Newton's second law, the vibration equation of the worktable is:

$$M \ddot{s}_u = -k_u [s_u - u(\lambda L, t)]. \quad (17)$$

Setting:

$$s_u(t) = S_u e^{i\omega_u t}. \quad (18)$$

By putting Eqs. (12), (14), (18) into the Eqs. (15-17), equations about  $D_1$ - $D_4$  and  $S_u$  can be got, and written as:

$$D \{D_1 \quad D_2 \quad D_3 \quad D_4 \quad S_u\}^T = 0. \quad (19)$$

For non-trivial solution, the coefficient matrix must be singular and one can get longitudinal frequency equation:

$$|D| = 0. \quad (20)$$

## 4. Calculation and analysis

### 4.1. Lateral vibration frequency

#### 4.1.1. The lateral vibration frequency of the beam without intermediate spring-mass system

If the worktable and the contact deformation between the worktable and the screw are not considered, the screw can be seen as Timoshenko beam, which supporting modal is shown in Fig. 3.

From Ref. [21], the first order lateral frequency of the classic Euler- Bernoulli beam under simply supported at both ends is:

$$\omega_{oss} = \left(\frac{\pi}{L}\right)^2 \sqrt{\frac{EI}{\rho A}}. \quad (21)$$

If  $\omega_{off}$  is the first order natural frequency of the Euler-Bernoulli beam under freely supported at both ends, and  $\omega_{ogg}$  is the first order natural frequency of the Euler-Bernoulli beam under clamed supported at both ends, the expression of both are:

$$\omega_{off} = \omega_{occ} = \left(\frac{4.73}{L}\right)^2 \sqrt{\frac{EI}{\rho A}} \tag{22}$$

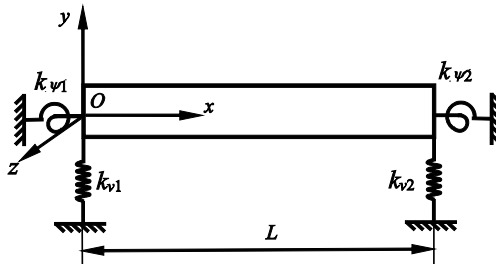


Fig. 3. The supporting model of the beam lateral vibration

If the lateral stiffness of bearing is enough large but the rotation stiffness around z axis is zero, the beam can be seen as simply supported; If the lateral stiffness and rotational stiffness of the bearing are close to zero, beam can be seen as freely supported; If the lateral stiffness and rotational stiffness of the bearing are both enough large, the beam can be clamped supported. Letting  $F = 0, \Omega = 0, k_v = 0, k_\psi = 0, M = 0$ , Eq. (11) reduces to  $4 \times 4$  orders. Letting  $E = 2.07 \times 10^{11}$  Pa,  $G = 8.3 \times 10^{10}$  Pa,  $\rho = 7.85 \times 10^3$  kg·m<sup>-3</sup>,  $L = 1000$  mm,  $d = 10$  mm, 20 mm, 40 mm respectively, the lateral vibration frequency is calculated. Letting  $k_{v1} = k_{v2} = 1000$  N·μm<sup>-1</sup>,  $k_{\psi 1} = k_{\psi 2} = 0$ , the frequency is written as  $\omega_{Tss}$ ; when  $k_{v1} = k_{v2} = 0, k_{\psi 1} = k_{\psi 2} = 0$ , the frequency is written as  $\omega_{Tff}$ ; when  $k_{v1} = k_{v2} = 1000$  N·μm<sup>-1</sup>,  $k_{\psi 1} = k_{\psi 2} = 1000 \times 10^6$  N·rad<sup>-1</sup>, the frequency is written as  $\omega_{Tcc}$ . The comparison between the result from Eq. (11) and Euler- Bernoulli beam is listed in Table 1.

Table 1. The frequency comparison between the Euler- Bernoulli beam and Timoshenko beam

	$\omega_{Oss}$	$\omega_{off} = \omega_{occ}$	$\omega_{Tss}$	$\omega_{Tff}$	$\omega_{Tcc}$
$d = 10$ mm	$1.2670 \times 10^2$	$2.8722 \times 10^2$	$1.2669 \times 10^2$	$2.1872 \times 10^2$	$2.8710 \times 10^2$
$d = 20$ mm	$2.5341 \times 10^2$	$5.7444 \times 10^2$	$2.5329 \times 10^2$	$4.3714 \times 10^2$	$5.7329 \times 10^2$
$d = 40$ mm	$5.0682 \times 10^2$	$1.1489 \times 10^3$	$5.0569 \times 10^2$	$8.7186 \times 10^2$	$1.1355 \times 10^3$

From Table 1, one can see that  $\omega_{Tss}, \omega_{Tff}$  and  $\omega_{Tcc}$  calculated from this paper, are close to the vibration frequency of the Euler-Bernoulli beam under simply, freely and clamed supported at both ends, respectively. But the vibration frequency of Timoshenko beam is slightly lower than that of the Euler-Bernoulli beam, moreover the larger the diameter is, the greater the influence of rotary inertia and shear deformation is, and the greater the difference between the both is.

The lateral vibration frequency of beam is calculated as a function of the radial supporting stiffness of the bearing which is shown in Fig. 4. From Fig. 4 one can see that the frequency varies when the radial supporting stiffness of the bearing increases. When  $k_{v1}$  and  $k_{v2}$  tend to infinity, the frequency is close to that under the simply supported.

Letting  $k_{v1} = k_{v2} = 1000$  N·μm<sup>-1</sup>, Fig. 5 is the first order lateral vibration frequency as function of the rotational stiffness of the bearing when  $d = 10$  mm,  $d = 20$  mm,  $d = 40$  mm respectively. The frequency increases with an increasing of the beam’s diameter due to the inertia. When  $k_{\psi 1} = k_{\psi 2} = 0$ , the frequency is close to that under simply supported and with rotational stiffness gradually increasing, the lateral vibration frequency increases to that under clamed supported ( $\omega_{Oss}$  and  $\omega_{Occ}$  are listed in Table 1).

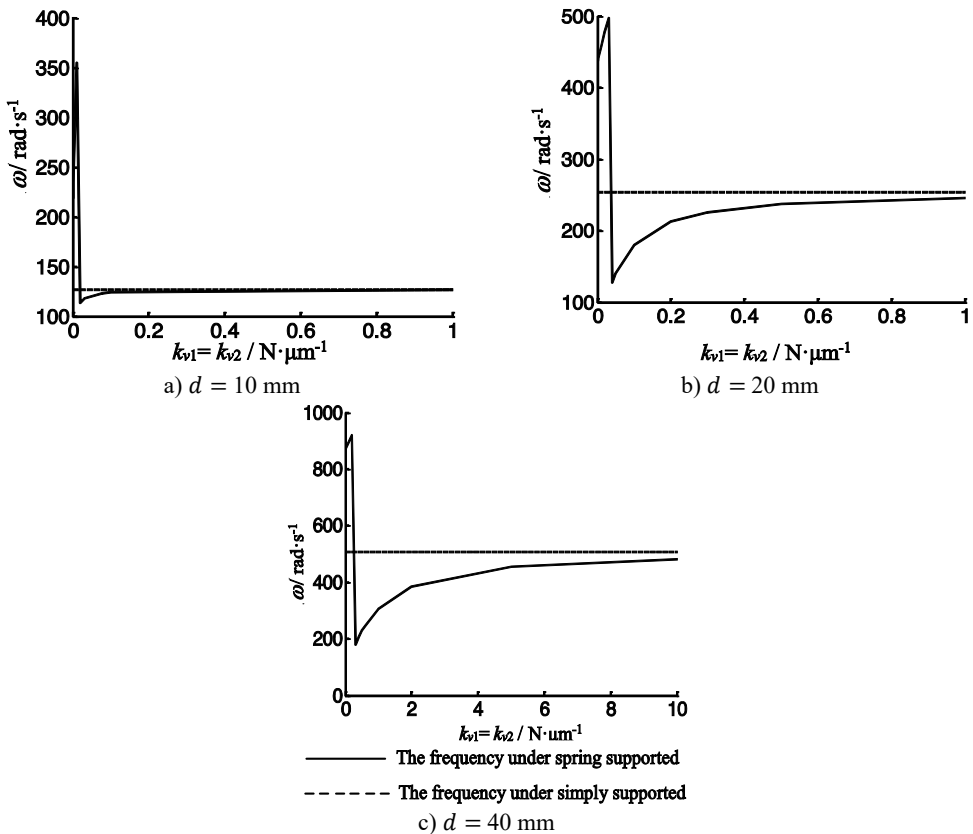


Fig. 4. The first order lateral frequency of the beam as function of the radial stiffness of the both bearing

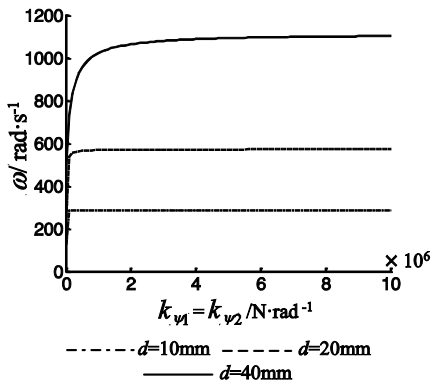


Fig. 5. The first order frequency of the beam as a function of the rotational stiffness of the both bearing

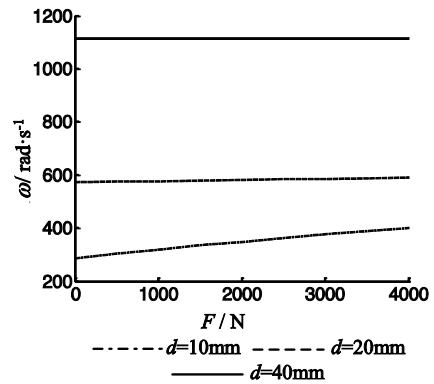


Fig. 6. The first order frequency of the beam as function of the pretension force

Fig. 6 is the first order lateral vibration frequency as function of the pre-tension force  $F$ . The vibration frequency increases with an increasing of the pre-tension force due to the increasing of the lateral stiffness. The smaller the diameter of the screw is the larger the effect of the pre-tension force is.

Fig. 7 is the first order lateral vibration frequency as function of the rotational speed of the beam. One can see that the vibration frequency increases with an increasing of rotational speed  $\Omega$  due to the gyroscopic effect. So, the larger the diameter of the screw is the larger the effect of rotational speed is.



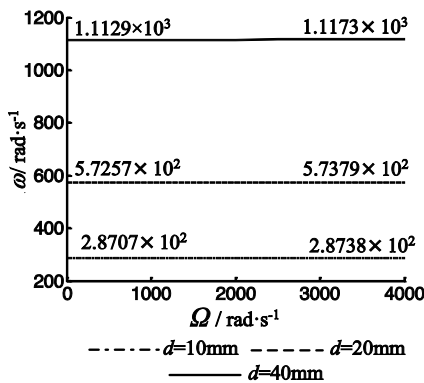
### 4.1.2. The lateral vibration frequency of the screw feeding system with intermediate spring-mass system

The lateral vibration of the screw feeding system is simulated by Ansys. The element type of the screw and the worktable are Beam188 and Mass21 respectively, and the contact stiffness of the lead screw and the worktable is neglected, and the boundary condition at the two ends is clamped. In Eq. (11), the stiffness of the bearing and the contact stiffness between the worktable and the screw are all  $200 \text{ N}\cdot\mu\text{m}^{-1}$ . When the worktable is at different position, the calculation results of Ansys are compared with the results of this paper written in Table 2. Because the stiffness of Ansys model is larger, the frequency of the system increases.

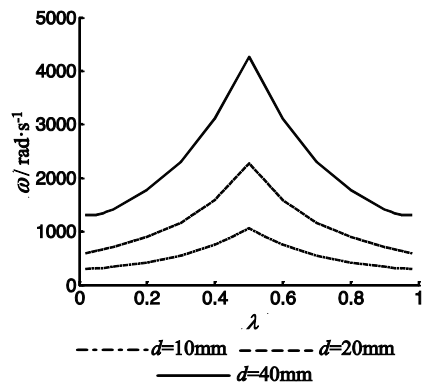
**Table 2.** The lateral vibration frequency comparison of the system between Ansys and this paper

	$\lambda = 0.25$	$\lambda = 0.5$	$\lambda = 0.75$
Ansys	1394	1633	1394
Eq. (11)	1162	1418	1162

Assuming the intermediate supporting stiffness of the screw  $k_v = 200 \text{ N}\cdot\mu\text{m}^{-1}$ ,  $k_\psi = 200 \times 10^6 \text{ N}\cdot\text{rad}^{-1}$ , the bearings' stiffness  $k_{v1} = k_{v2} = 200 \text{ N}\cdot\mu\text{m}^{-1}$ ,  $k_{\psi1} = k_{\psi2} = 200 \times 10^6 \text{ N}\cdot\text{rad}^{-1}$ , when the worktable moves along the screw, the lateral vibration frequencies of the screw feeding system are shown in Fig. 8 according Eq. (11). From Fig. 8, one can see that the vibration frequency of the system increases as its diameter increases, and the frequency varies when the worktable moves. When the worktable moves along the screw from the left end to the right end, the frequency is symmetrical about the midpoint of the screw. When the worktable is at middle of the screw, the lateral frequency is biggest. The vibration frequencies of the system with different bearings' stiffness are shown in Fig. 9. When the two bearings' stiffness are not same, the frequency is not symmetrical, and the larger the diameter of the screw is, the more obvious asymmetry of the frequency is.



**Fig. 7.** The first order frequency of the beam as function of the rotational speed



**Fig. 8.** The vibration frequency of the feeding systems a function of position coefficient of the worktable

Letting  $\lambda = 0.5$ , Fig. 10 shows the vibration frequency of the system depends on intermediate stiffness. From Fig. 10, one can see that when the intermediate stiffness increases, the frequency increases. Fig. 11 shows that the frequency of the system depends on the mass of the worktable, and Fig. 12 shows the frequency of the system depends on the contact stiffness between the worktable and the guide way. One can see that the frequency decreases with the mass increasing, and increases with the contact stiffness increasing between the worktable and the guide way.

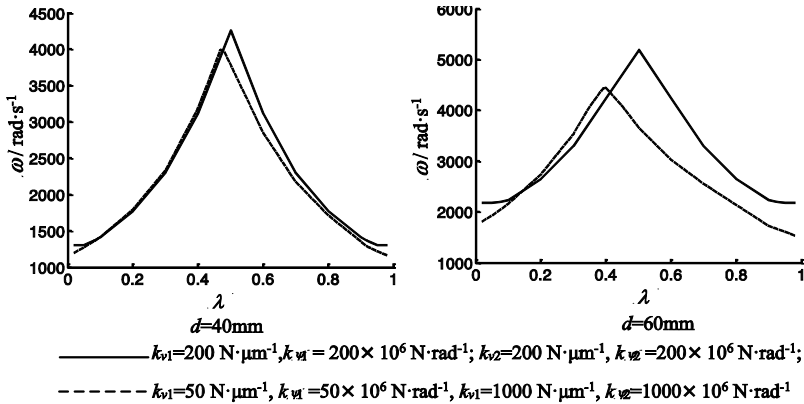


Fig. 9. The vibration frequency of the feeding system as a function of the bearings stiffness

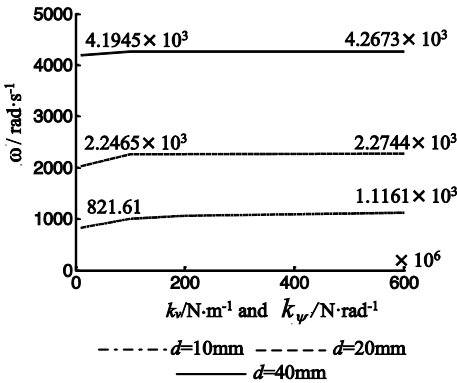


Fig. 10. The vibration frequency of the feeding system as function of the intermediate support stiffness

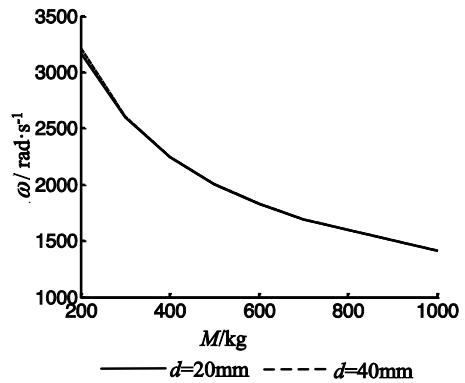


Fig. 11. The vibration frequency of the feeding system as function of the worktable mass

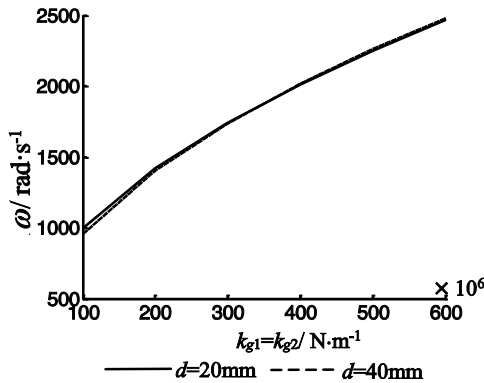


Fig. 12. The vibration frequency of the feeding system as function of contact stiffness between the worktable and the guide way

## 4.2. The longitudinal vibration frequency

### 4.2.1. The longitudinal vibration frequency of the rod without the intermediate spring-mass system

Without considering the worktable, the longitudinal vibration of the screw system can be simplified as a rod with spring supported which is shown in Fig. 13. The bearing's supporting of

the rod's both ends can be simplified as two kinds, one is clamped when the bearing's axial stiffness is big enough, and the other is free when the bearing's axial stiffness is small enough.

When the supporting of rod's both ends is clamped or free, the longitudinal vibration frequency equation all are:

$$\omega_{ucc} = \omega_{uff} = \frac{r\pi}{L} \sqrt{\frac{E}{\rho}}, \quad r = 1, 2, \dots, n, \quad (23)$$

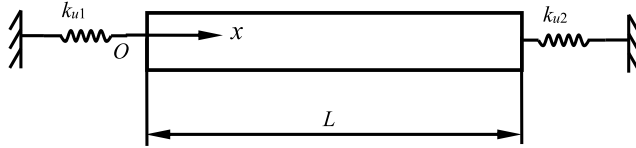


Fig. 13. The longitudinal vibration supporting model of the rod

In this paper, letting  $k_u = 0$ ,  $M = 0$ , Eq. (20) reduces to  $2 \times 2$  orders, which is the longitudinal frequency equation of the rod under spring supported at both ends. Letting  $E = 2.07 \times 10^{11}$  Pa,  $G = 8.3 \times 10^{10}$  Pa,  $\rho = 7.85 \times 10^3$  kg·m<sup>-3</sup>,  $L = 1000$  mm, the frequency of longitudinal vibration is calculated. Letting the bearing's supporting stiffness of the screw both ends is same, and the diameter  $d = 10$  mm, 20 mm, 40 mm respectively, the first order longitudinal vibration frequency with the axial supporting stiffness of both ends bearings is shown in Fig. 14. When bearing stiffness of both ends is 0, the vibration frequency is under freely supported at both ends. With the increasing of bearing stiffness, longitudinal vibration frequency of the rod increases first, but falls sharply at a certain critical stiffness  $k_{uc}$ , and then gradually increases to reach the natural frequency under clamped supported at both ends. Fig. 14 shows that the bearing supporting can be simplified as freely supported at both ends if the bearing stiffness is smaller than the critical stiffness  $k_{uc}$ , and can be simplified as clamped supported at both ends if the bearing stiffness is more than the critical stiffness  $k_{uc}$ . The critical stiffness of  $k_{uc}$  is related to the diameter of the rod and increases with the rod diameter increasing. The bearing supporting stiffness on both ends should be reasonable according to the diameter of the rod to avoid the vibration instability.

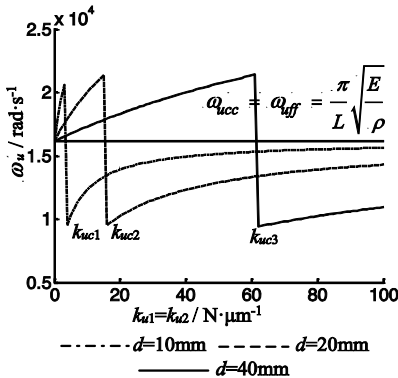


Fig. 14. The vibration frequency of the rod as function of the both bearing stiffness

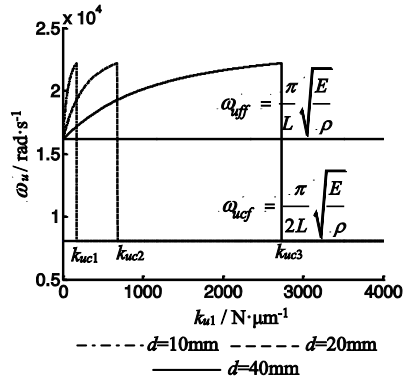


Fig. 15. The vibration frequency of the rod as function of the bearing stiffness when  $k_{u2} = 0$

When the rod' length or the temperature variety is larger, it is often used to be under clamped supported at one end and freely supported at the other end, in which the vibration frequency equation is:

$$\omega_{ucf} = \frac{(2r - 1)\pi}{L} \sqrt{\frac{E}{\rho}}, \quad r = 1, 2, \dots, n. \quad (24)$$

Letting  $k_{u2} = 0$ , the vibration frequency with the axial supporting stiffness of one end bearings is calculated from this paper shown in Fig. 15. When  $k_{u1} = 0$ , the vibration frequency is under freely supported of both ends. With the increasing of  $k_{u1}$  the vibration frequency increases gradually, and falls sharply to the frequency under clamped supported at one end at a critical stiffness  $k_{uc}$ . From Fig. 15, one can see that if the screw is under freely supported at one end, only when the other end supporting stiffness is greater than the critical stiffness  $k_{uc}$ , it can be seen as clamped supporting.

From Eqs. (23-24) one can see that, if the stiffness of the bearing is not considered, the longitudinal vibration frequency of the rod does not change with the diameter of the screw. But from Fig. 14-15 one can see that, if considering the stiffness of the bearing, the frequency of the longitudinal vibration of the rod varies with the diameter.

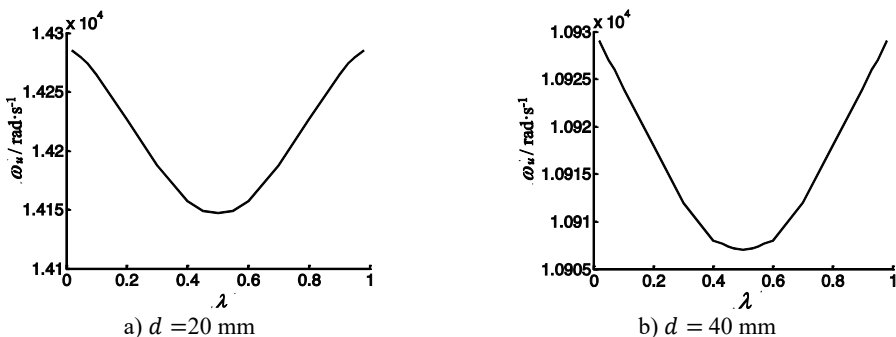
### 4.2.2. The longitudinal vibration frequency of the screw feeding system with the intermediate spring-mass system

Using Ansys to simulate the longitudinal vibration of the screw feeding system, the element types of the screw and the worktable are Link180 and Mass21 respectively, and the contact stiffness of the screw and the worktable is neglected, and the boundary condition at the two ends is clamped. According Eq. (20), the stiffness of the bearing is  $1000 \text{ N}\cdot\mu\text{m}^{-1}$ , and the contact stiffness between the screw and the worktable is  $300 \text{ N}\cdot\mu\text{m}^{-1}$ . The calculation results of ANSYS and this paper are shown in Table 3.

**Table 3.** The longitudinal vibration frequency comparison of the system between Ansys and this paper

	$\lambda = 0.25$	$\lambda = 0.5$	$\lambda = 0.75$
Ansys	$1.2526 \times 10^4$	$1.2466 \times 10^4$	$1.2526 \times 10^4$
Eq. (20)	$1.4207 \times 10^4$	$1.4147 \times 10^4$	$1.4207 \times 10^4$

Considering the worktable moving and the contact deformation between the worktable and the screw, the longitudinal vibration frequency is calculated according the Eq. (20). Letting  $k_{u1} = k_{u2} = 1000 \text{ N}\cdot\mu\text{m}^{-1} > k_{uc}$ , the frequency depends on the position coefficient  $\lambda$  is shown in Fig. 16. The frequency varies when the worktable moves along the screw. When the worktable moves along the screw from the left end to the right end, the frequency is symmetrical about the midpoint of the screw. When the worktable is at the middle of the screw, the frequency is lowest.



**Fig. 16.** The vibration frequency of the feeding system as a function of position coefficient of the worktable under clamped-clamped supported

If the screw is long or the temperature varies large, the screw is often under clamped-freely

supported. So, letting  $k_{u1} = 3000 \text{ N}\cdot\mu\text{m}^{-1}$  and  $k_{u2} = 0$ , the frequency depends on the position coefficient  $\lambda$  is shown in Fig. 17. The vibration frequency decreases when the worktable moves along the screw from the left end to the right end.

Fig. 18 shows that the longitudinal vibration frequency of the system depends on the mass of the worktable, and the frequency decreases with the mass increasing. Fig. 19 shows the frequency of the system depends on the contact stiffness between the worktable and the screw. The frequency increases with the contact stiffness increasing, but when the contact stiffness is large enough, the frequency almost remains unchanged with the contact stiffness increasing.

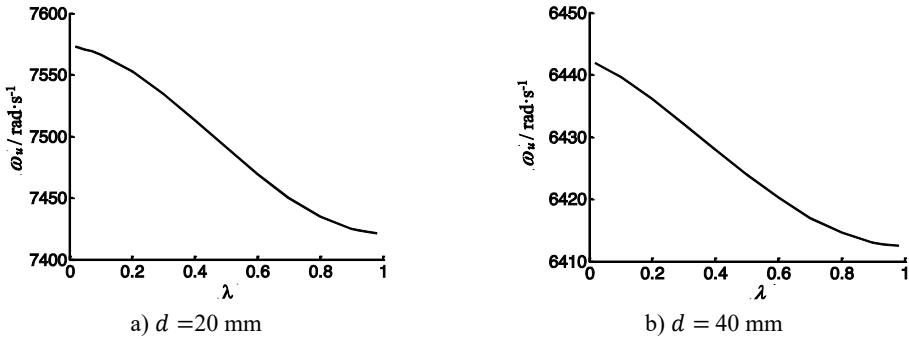


Fig. 17. The vibration frequency of the screw feeding system as a function of position coefficient of the worktable under clamed-free supported

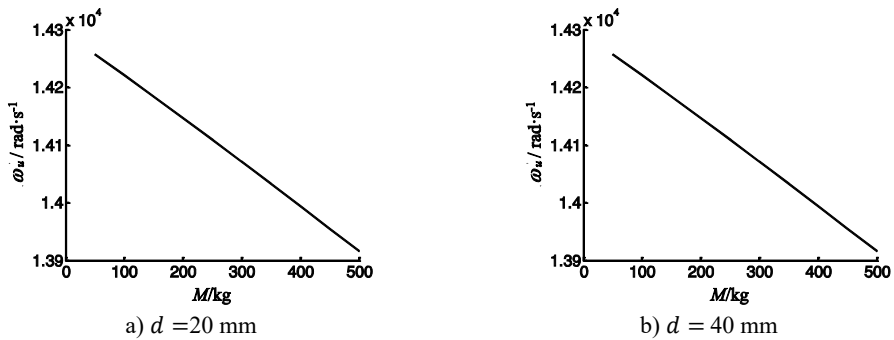


Fig. 18. The vibration frequency as function of the worktable mass

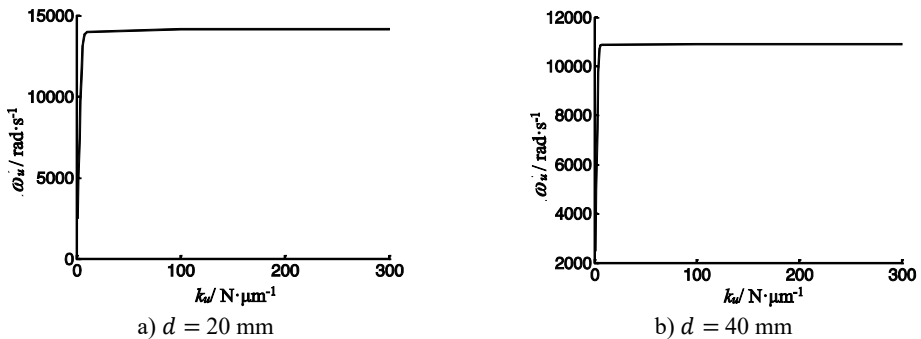


Fig. 19. The vibration frequency as function of the contact stiffness between the screw and the worktable

5. Conclusions

Through the above analysis, the following conclusions are obtained:

- 1) The diameter, the pre-tension force and the rotational speed have effect on the lateral

vibration frequency of the beam, and when the diameter, the pre-tension force or the rotating speed increases, the frequency increases.

2) The lateral vibration frequency varies when the worktable moves along the screw from the left end to the right end, and the frequency is biggest when the worktable is at the middle of the screw. If the bearings stiffness of the both ends is same, the frequency is symmetrical about the midpoint of the screw; and if the bearings stiffness is not same, the frequency is not symmetrical. The larger the diameter of the screw is, the more obvious asymmetry of the frequency is.

3) The longitudinal vibration frequency of the rod varies with the bearings stiffness, and there is a critical stiffness  $k_{uc}$ . The critical stiffness is related to the diameter of the rod and increases with the rod diameter increasing. The bearing supporting stiffness of both ends should be reasonable in order to avoid the vibration instability.

4) The longitudinal frequency of the screw feeding system varies when the worktable moves along the screw. When the screw is under clamed-clamed supported, the frequency is symmetrical about the midpoint of the screw. When the worktable is at the middle of the screw, the frequency is lowest. But when the screw is under clamed-freely supported, the longitudinal vibration frequency decreases when the worktable moves along the screw from the left end to the right end.

5) The mass of the worktable and the contact stiffness of the system have effect on the longitudinal frequency of the system. The frequency decreases with mass of the worktable increasing and increases with the contact stiffness increasing.

## Acknowledgement

The author would like to thank the Henan Polytechnic University for financially supporting the part of the work under Contract B2011-92.

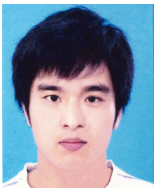
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